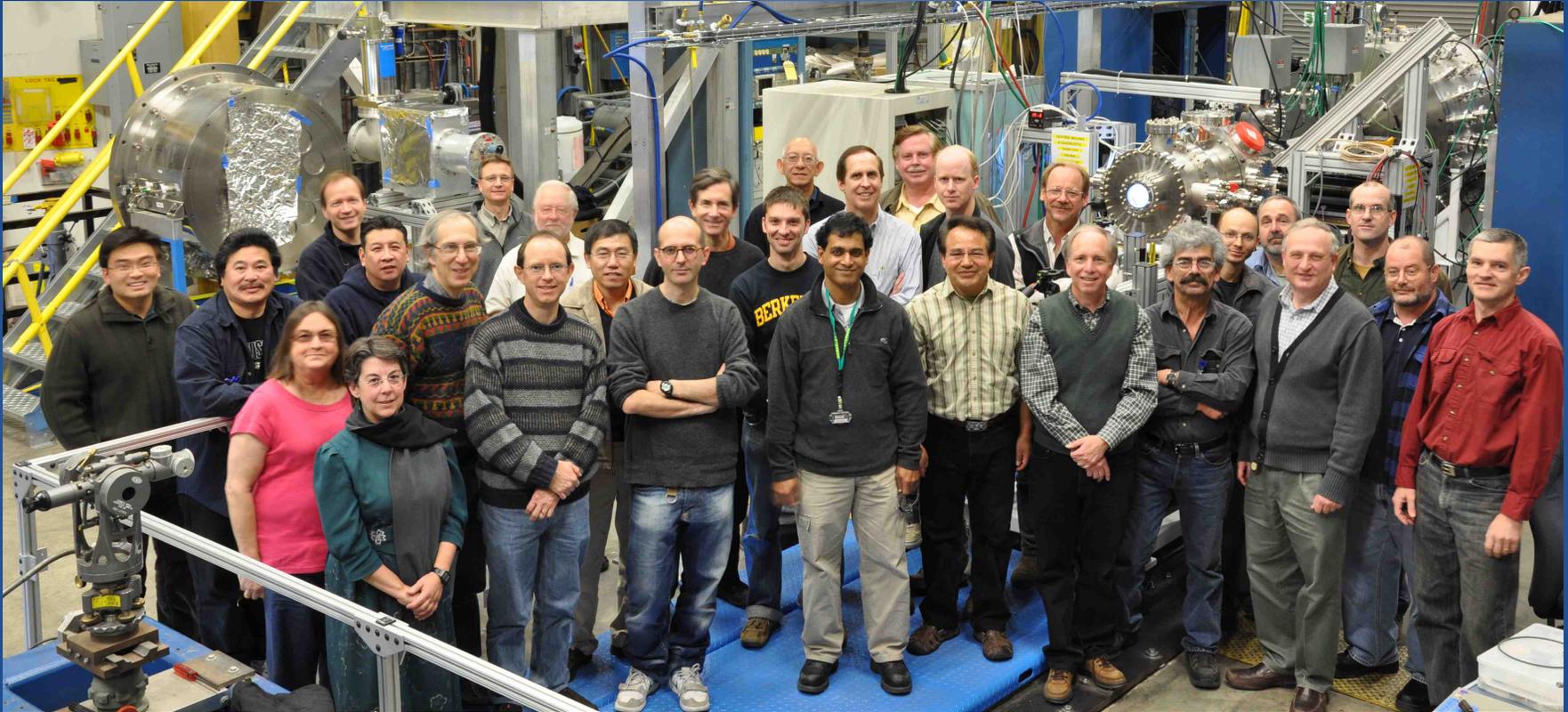


Manipulating Ion Beams for High Energy Density Physics Studies

Steve Lidia (HIFS-IBT, LBNL)
Center for Beam Physics Seminar
29 January 2010

HIFS-VNL Team



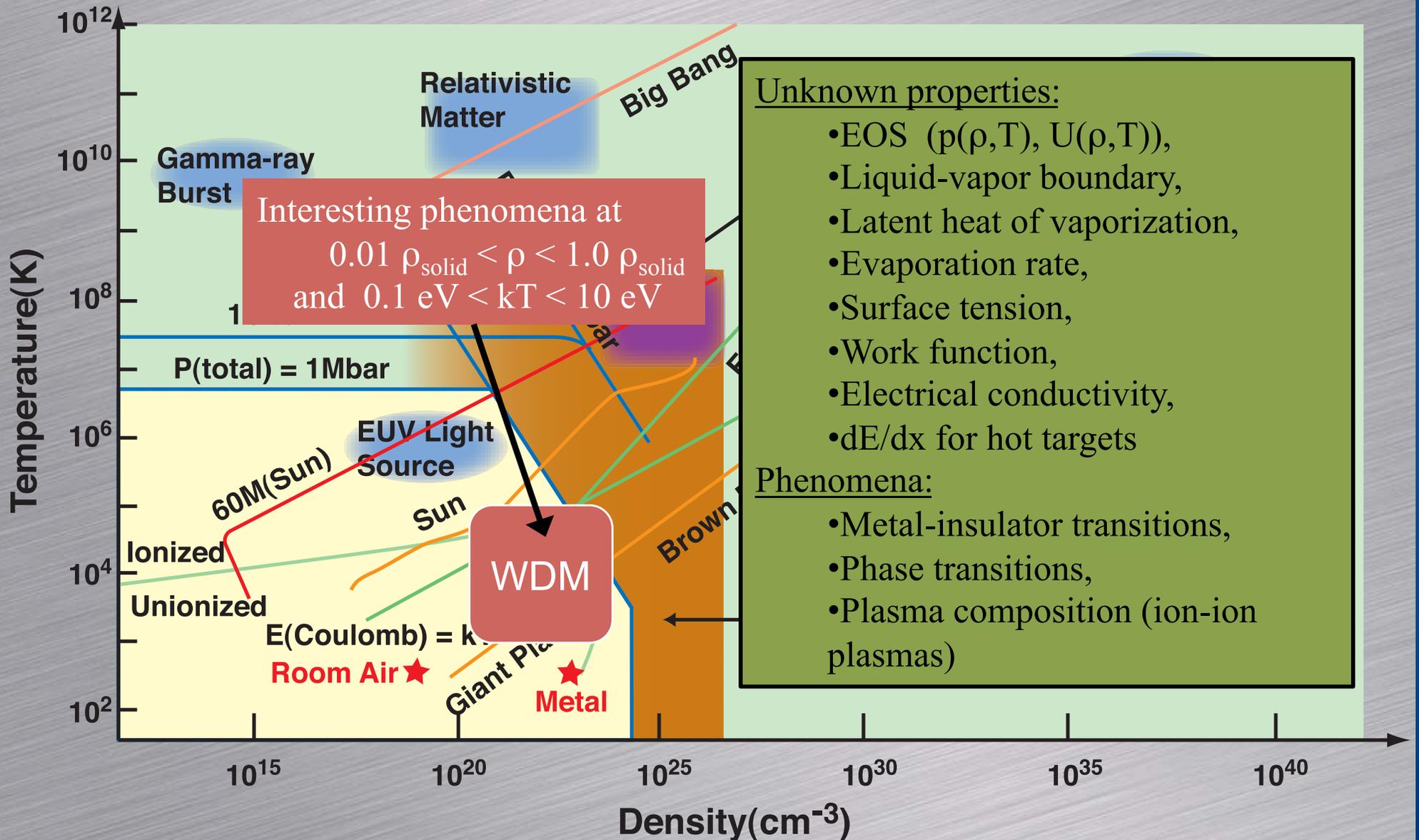
Not pictured:

G. Bazouin, F. Becker, J. Coleman, R. Davidson, M. Dorf, E. Gilson,
I. Kaganovich, M. Kauffeldt, D. Ogata, A. Sefkow, K. Van den Bogert,
J-L Vay, D. Welch, C. Wootton, S. Yu

Introduction

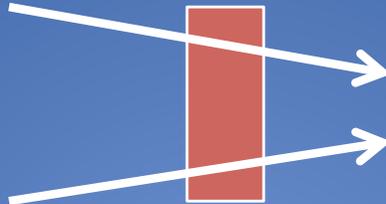
- High Energy Density (HED) and Warm Dense Matter (WDM)
- Why Ion Beams?
- Manipulating Ion Beams for HEDP
 - Stage 1: Neutralized transport and focus (NTX)
 - Stage 2: Neutralized drift compression (NDCX-I)
 - Stage 3: Add acceleration schedule (NDCX-II)
 - Stage 4: Pulse shaping and shock timing (IB-HEDPX)

HEDP Space



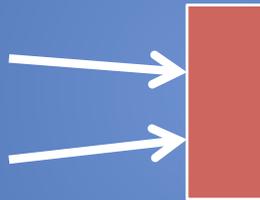
What's so interesting about ions?

Foils heated by hard x-rays



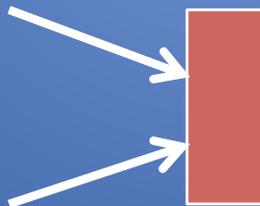
XFEL heating ->
small energy/volumes
($\sim 1\mu\text{m}$ radius, $\sim 0.1\mu\text{m}$ thickness)
May limit diagnostic accuracy.
Poor coupling to lower Z.

Supersonically heated foams or
low Z material (thermal x-rays)



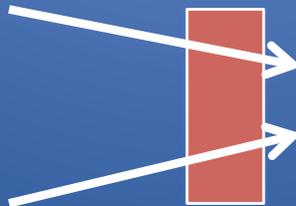
MJ of soft x-rays available on Z
machine but limited number of
shots.

Shock compressed and heated
thin foils



Lasers absorb at $\rho_{\text{critical}} \ll \rho_{\text{solid}}$
-> large density/pressure gradients

Tailored Bragg-peak ion beams
heating thin foils



Fast heating of a solid with
penetrating ions
-> more volume/lower gradients
-> more accurate EOS

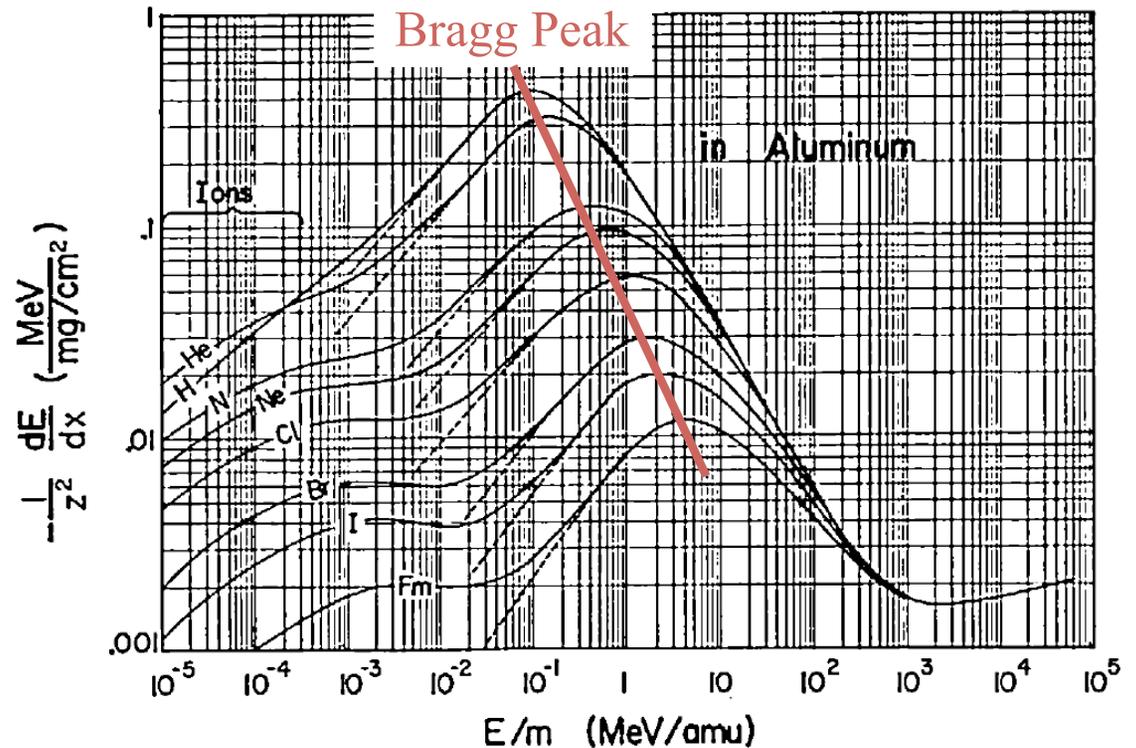
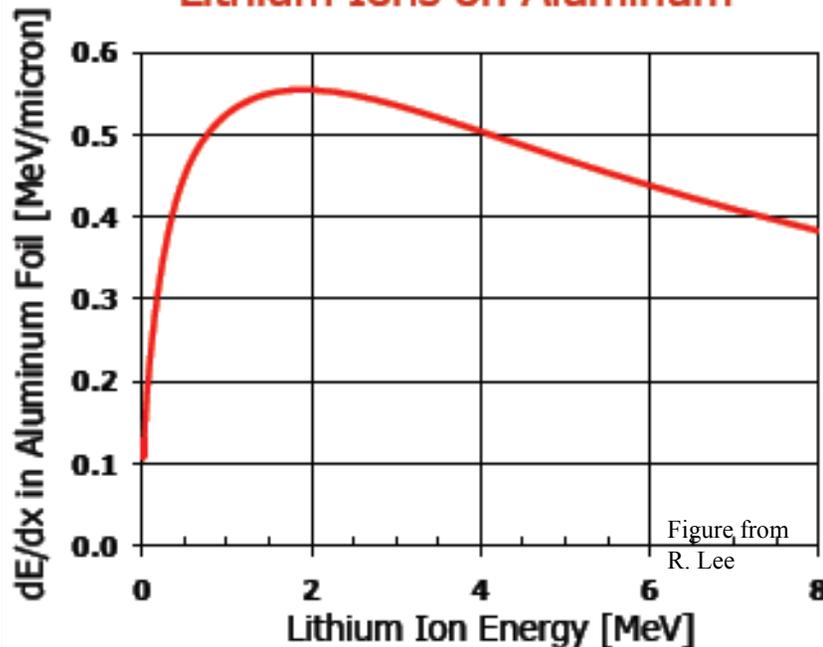
Ion Beams on Solids

Range matched to energy and species,

For < 1 MeV, can use lighter ions, e.g. He^+ ,

For > 10 MeV, can use heavier ions, e.g. C^+

Lithium Ions on Aluminum



Ion beam source limitations

- Multicusp Plasma, Pulsed ECR
 - High residual gas pressure – results in resonance charge exchange processes that increase energy spread
 - High ion temp ($\sim 2\text{eV}$), high current density (100 mA/cm^2)
- Contact Ionization/Aluminosilicate
 - Produces alkali metal ions (with low ionization potential) from a solid surface with high work function (e.g. tungsten) at high temperature.
 - Ion temperature $\leq 0.3\text{ eV}$; $J \leq 20\text{ mA/cm}^2$.
 - It had been used widely for ion thrusters in the early space-rocket program.
- Metal Vapor Vacuum Arc (MEVVA) sources
 - $J \sim 10\text{'s mA/cm}^2$
 - Low pressure, no residual gas
 - Mixed ion charge states

High current / High brightness beams

$$B \equiv \left(\frac{2}{\pi^2} \right) \frac{I}{\varepsilon^2}$$

At the ion source:

$$\varepsilon \equiv 2r \sqrt{\frac{kT_i}{mc^2}}$$

thus

$$B \propto \frac{J}{T}$$

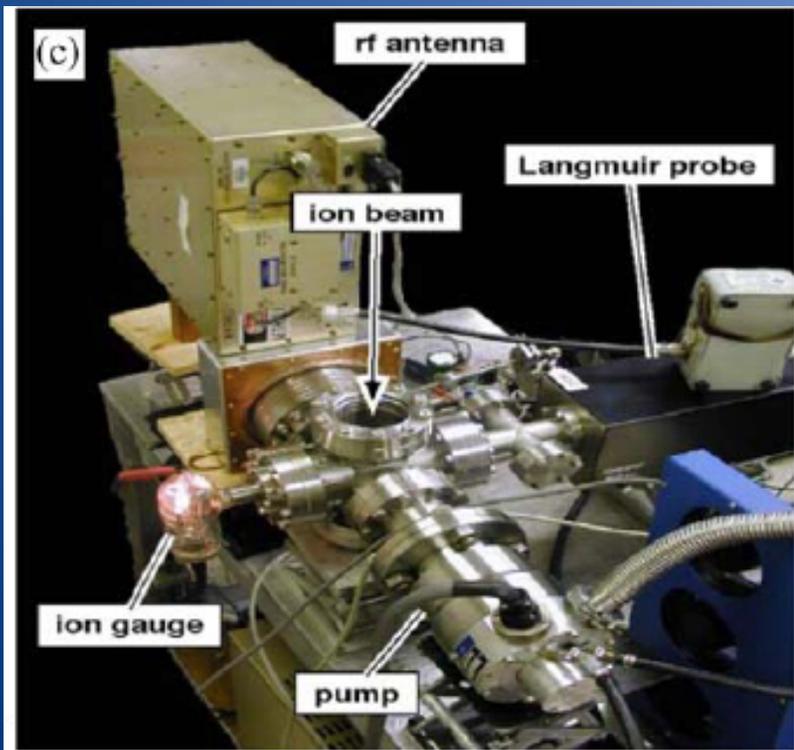
- High brightness means either low emittance, or high current, or both.
- At the ion source, high brightness is obtained by either high current density, or low ion temperature, or both.
- Achieving both high current and high brightness together is difficult at the ion source.
- Typical ions are (K, Li, Na,) Bi, Pb, Hg, Cs, Xe, Rb, K, Ar, and Ne.

Neutralization boosts final intensity

- Achieving mm to sub-mm spot sizes with high perveance, high brightness beams requires
 - Excessively large stored energy in final focus solenoids, or
 - Removing repulsive space charge forces and using reasonably sized focusing magnets
- Neutralized plasma requirements
 - $Z_b N_e > N_b$
 - Denser electron plasma densities closer to target plane

Plasma Technology

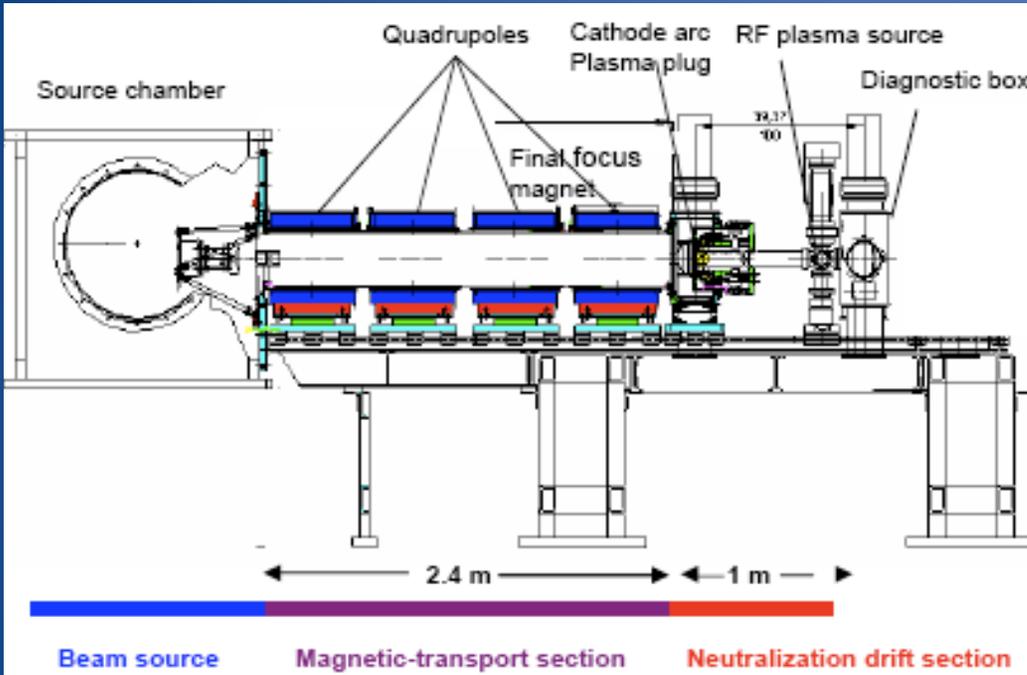
- Ferroelectric plasma source (FEPS)
 - ($N_e \sim 10^{10} / \text{cm}^3$)
- RF sources
 - ($N_e \sim 10^{11} / \text{cm}^3$)
- Cathodic Arc plasma source (CAPS)
 - ($N_e \sim 10^{14} / \text{cm}^3$)



Introduction

- HEDP and WDM
- Why Ion Beams?
- Manipulating Ion Beams for HEDP
 - Stage 1: Neutralized transport and focus (NTX)
 - Stage 2: Neutralized drift compression (NDCX-I)
 - Stage 3: Add acceleration schedule (NDCX-II)
 - Stage 4: Pulse shaping and shock timing (IB-HEDPX)

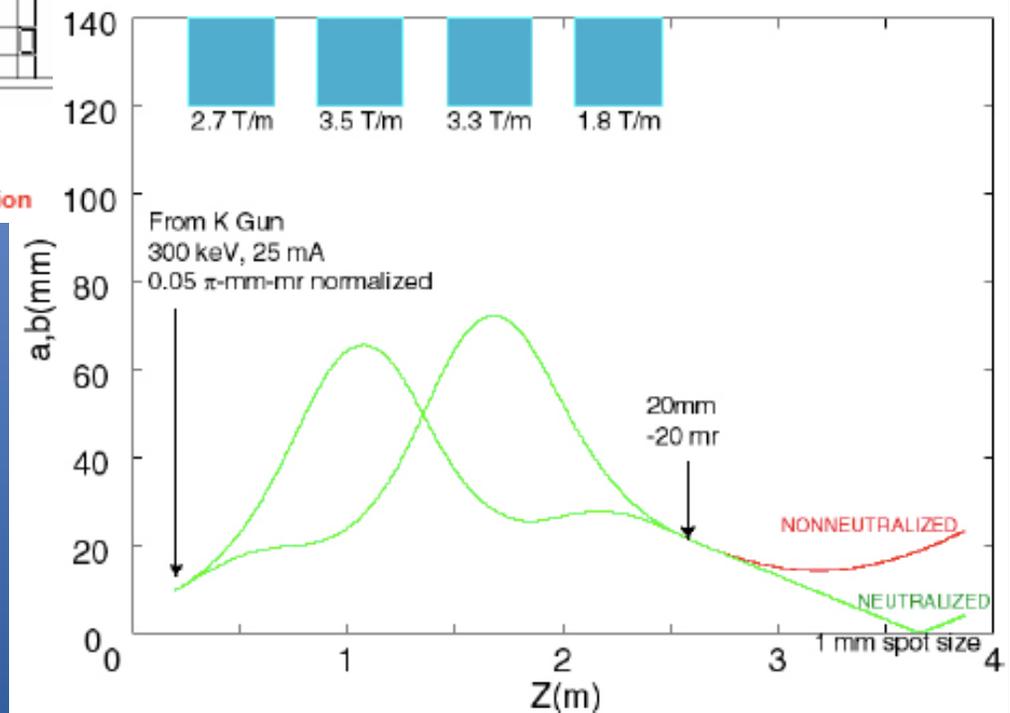
NTX: Neutralized Transport Experiment



Introduce neutralizing plasmas in final focus beamline

NTX Parameters:

- 300kV, 25mA, 10 μ sec K⁺ beam
- 4 transport magnetic quadrupoles
- 2 MEVVA density $\sim 10^{10}$ - 10^{11} cm⁻³
- RF plasma density $\sim 10^{11}$ - 10^{12} cm⁻³

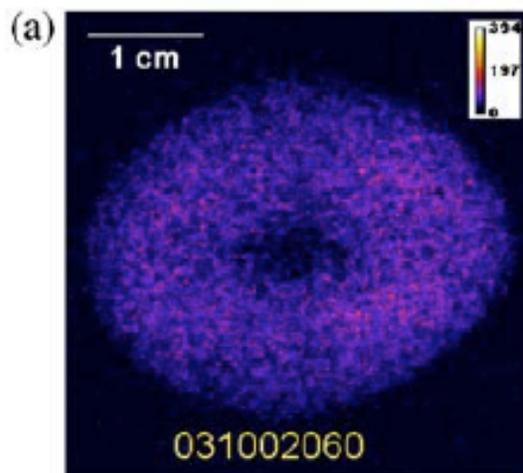


Neutralized Transport Results

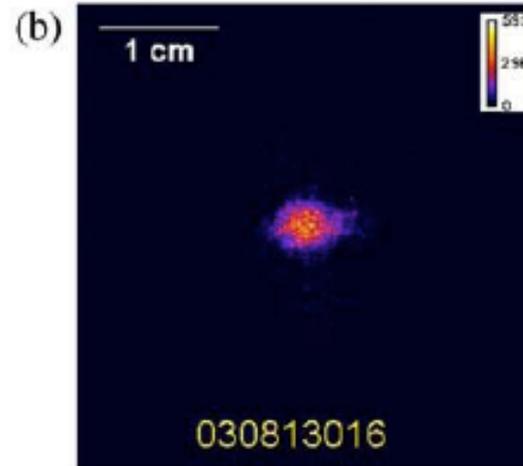
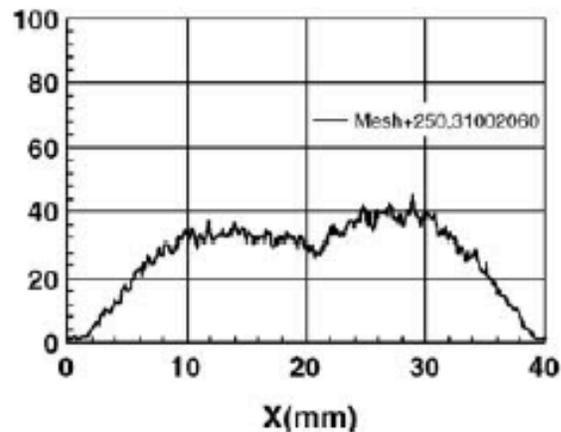
No neutralization

MEVVA only

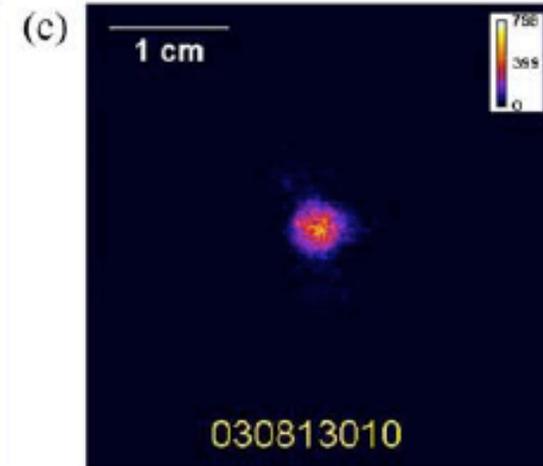
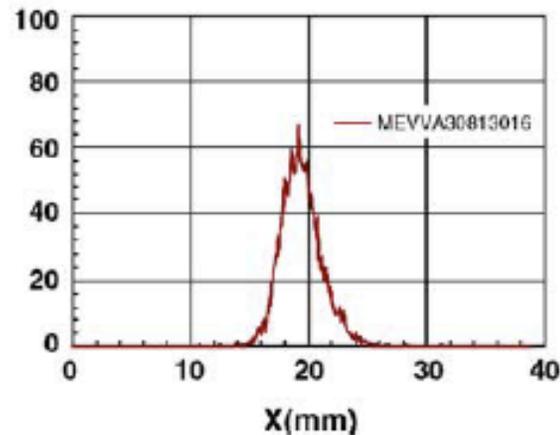
MEVVA + RF plasma



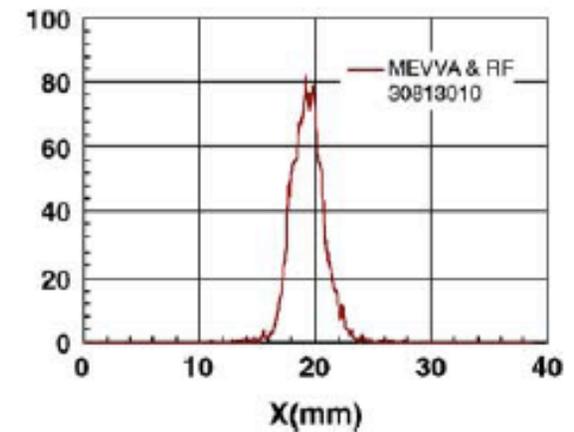
FWHM: 2.71 cm



FWHM: 2.83 mm



FWHM: 2.14 mm

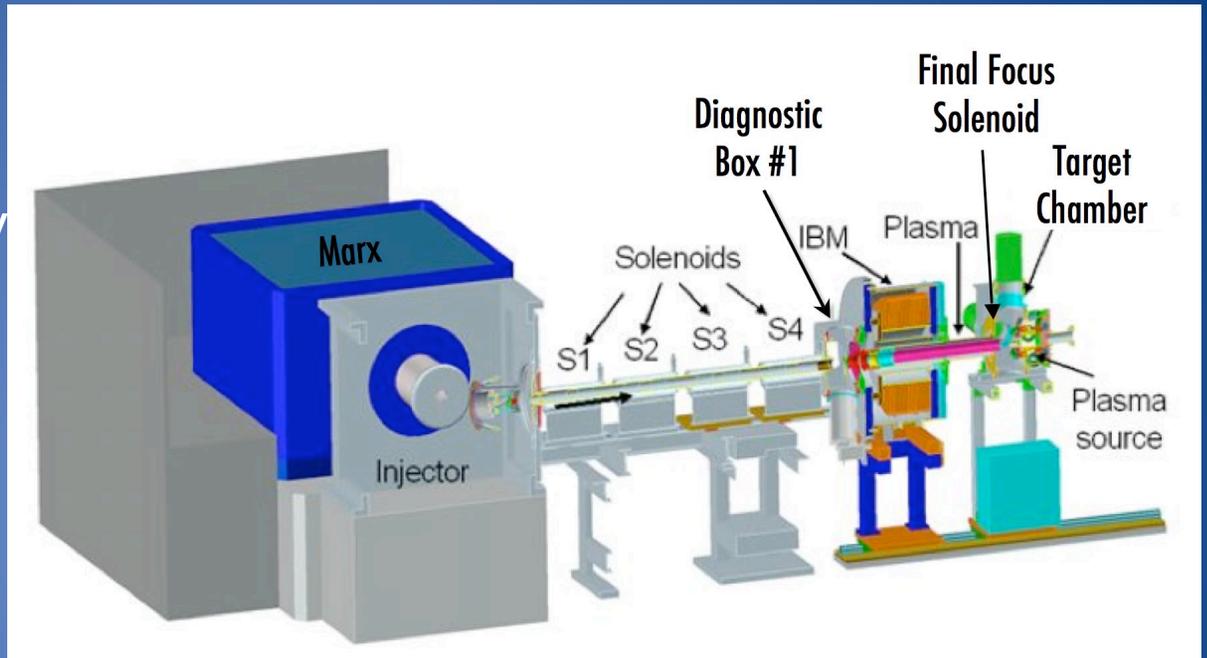


Introduction

- HEDP and WDM
- Why Ion Beams?
- Manipulating Ion Beams for HEDP
 - Stage 1: Neutralized transport and focus (NTX)
 - Stage 2: Neutralized drift compression (NDCX-I)
 - Stage 3: Add acceleration schedule (NDCX-II)
 - Stage 4: Pulse shaping and shock timing (IB-HEDPX)

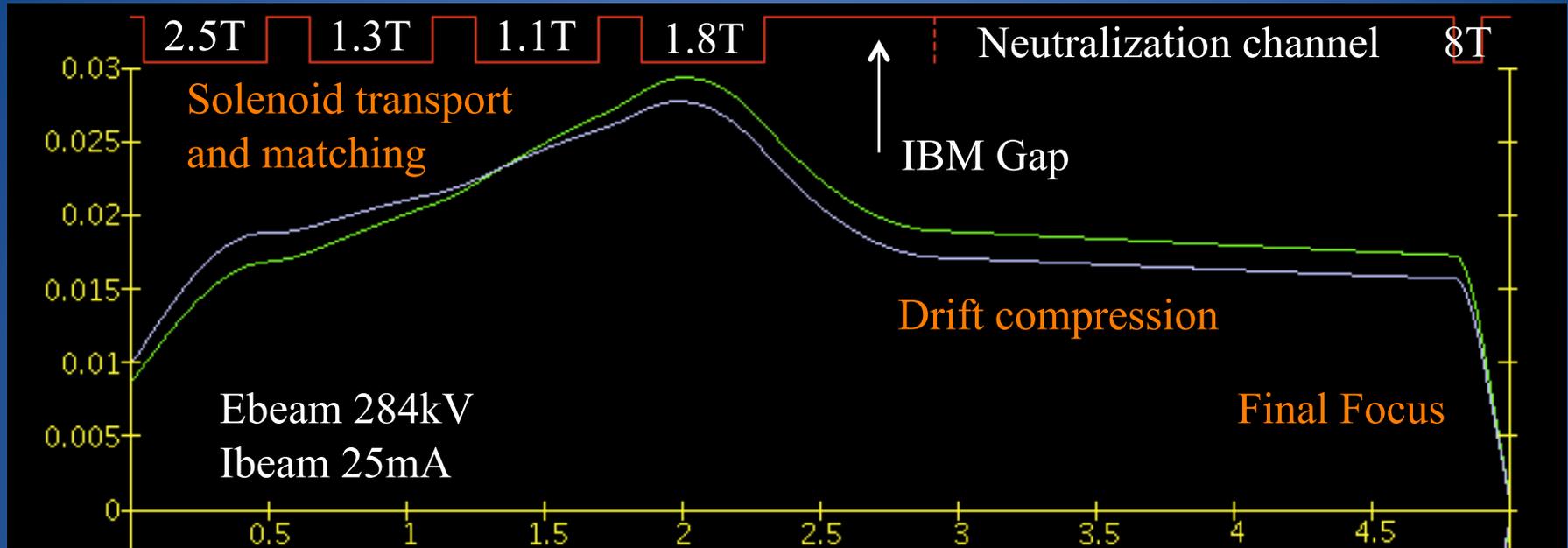
NDCX-I

Beam species K^+
 Beam current 25 - 40 mA
 Beam energy 270 - 350 kV
 Beam power ~7-14 kW
 Energy spread ~0.2 keV
 Pulse length ~3-20 μ sec



Injector solenoids	~3T peak field
Final focus solenoid	8T peak field
Neutralizing plasmas	
FEPS	$2 - 8 \cdot 10^{10} / \text{cm}^3$
4 MEVVA/FCAPS	$0.9 - 6 \cdot 10^{13} / \text{cm}^3$
Compressing Voltage Swing	~150 keV / 500 nsec

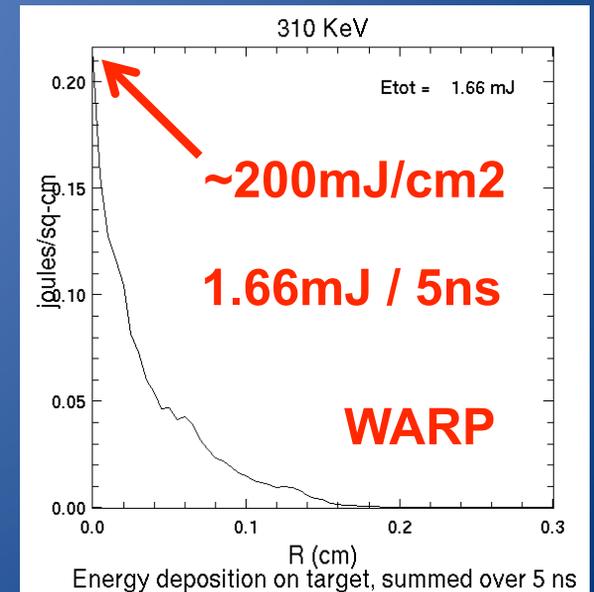
Envelope Dynamics



Uncompressed Beam

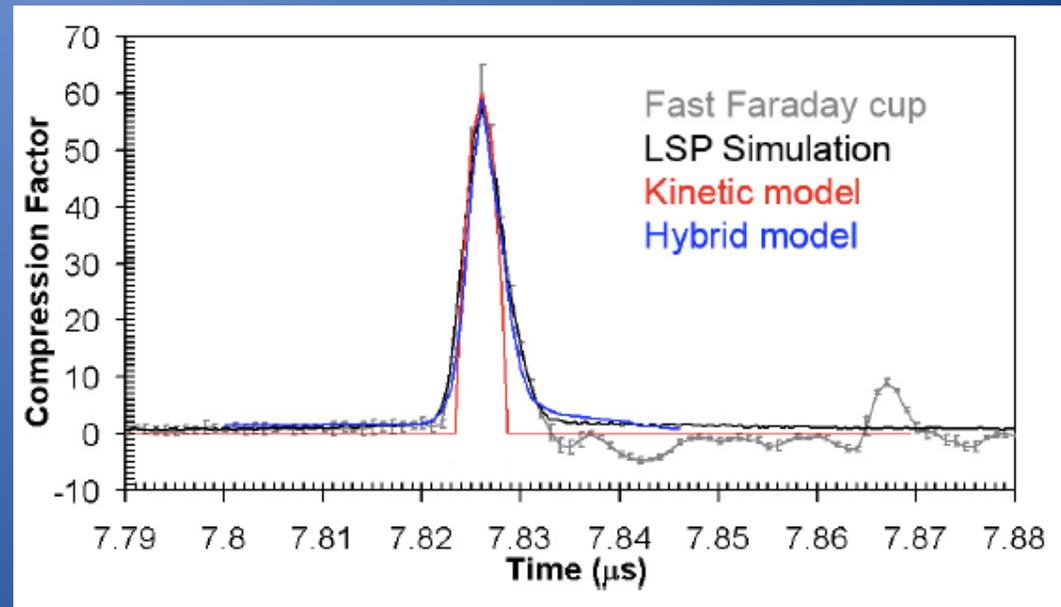
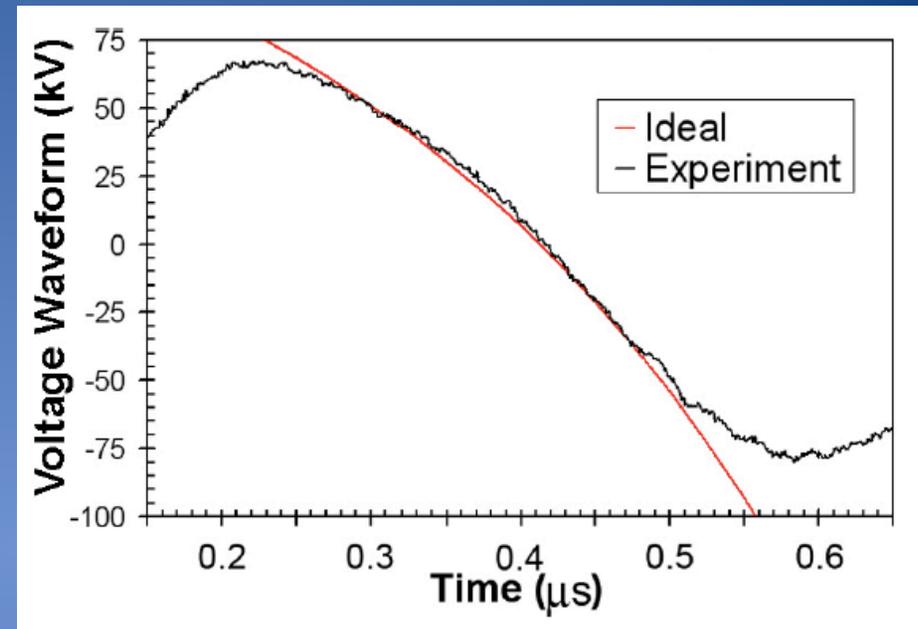
Beam compresses ballistically in the drift compression stage.

Compressed pulse fluence at target plane



Drift Compression

- Simple ballistic compression of non-relativistic beams.
- Use a programmable pulser to impart time-dependent velocity ramp.
- Neutralize beam space charge to improve compression factor.
- Trade some volt-seconds to add acceleration.



NDCX-I Emittance limits

- Transverse ($\beta_z \sim 0.004$)
 - Gun exit: $4\epsilon_{nrms} = 0.088 \pi \text{ mm mrad} \sim 2 (2\pi r_{emitter} v_{Tx})$
 - IBM entrance: $4\epsilon_{nrms} = 0.088 \pi \text{ mm mrad}$
- Longitudinal
 - $T_z = (2.4 \pm 1.6) 10^{-2} \text{ eV}$ K+ beam itself
 - $T_z = (4.7 \pm 2.8) 10^{-2} \text{ eV}$ K+ beam after neutralization

- Final Focus Optics

$\sim 25\text{mrad}$ convergence angle – $r_{spot} \sim 0.9\text{mm}$

$$r_{spot} \cong \frac{\epsilon}{r'_{conv}}$$

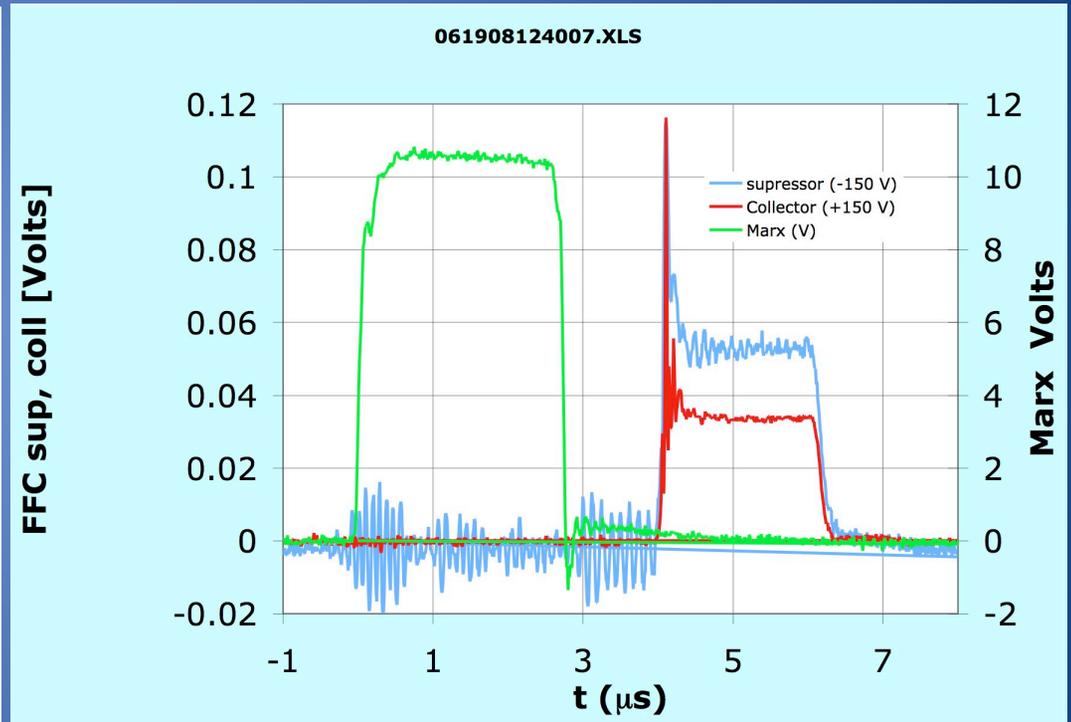
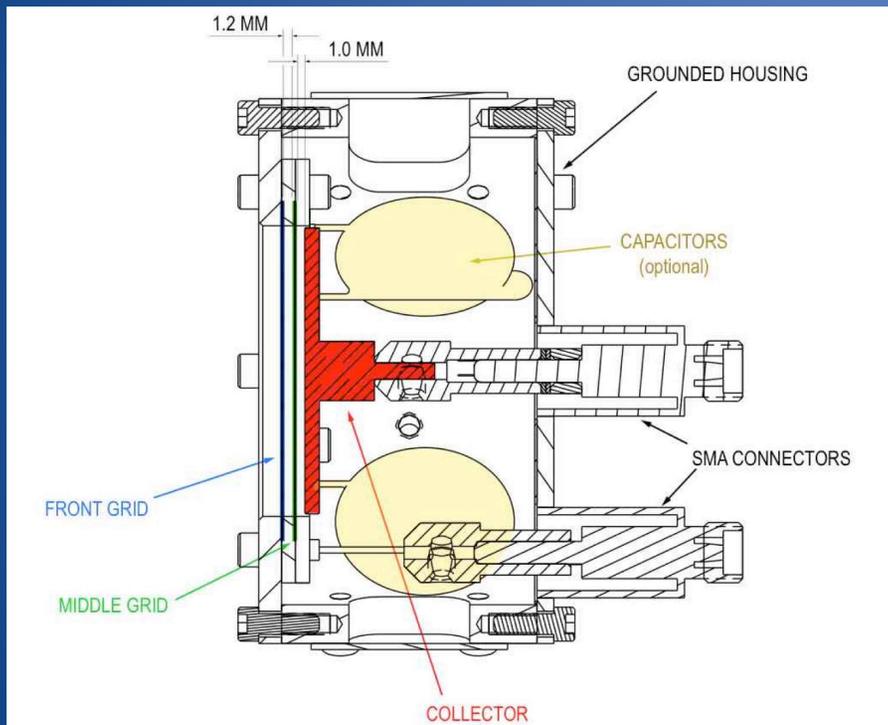
- Longitudinal Compression

$\sim 100\text{-}200\text{ps}$ for NDCX-I ($L_{drift} \sim 1.5 - 3 \text{ m}$)

$$t_{min} = \frac{L_{drift}}{2c} \left[\frac{m_i c^2}{E_i} \right] \sqrt{\frac{T_i}{m_i c^2}}$$

Target Chamber Beam Diagnostics

- **Fast Faraday Cup** ~ 1 ns response time (ion transit time)

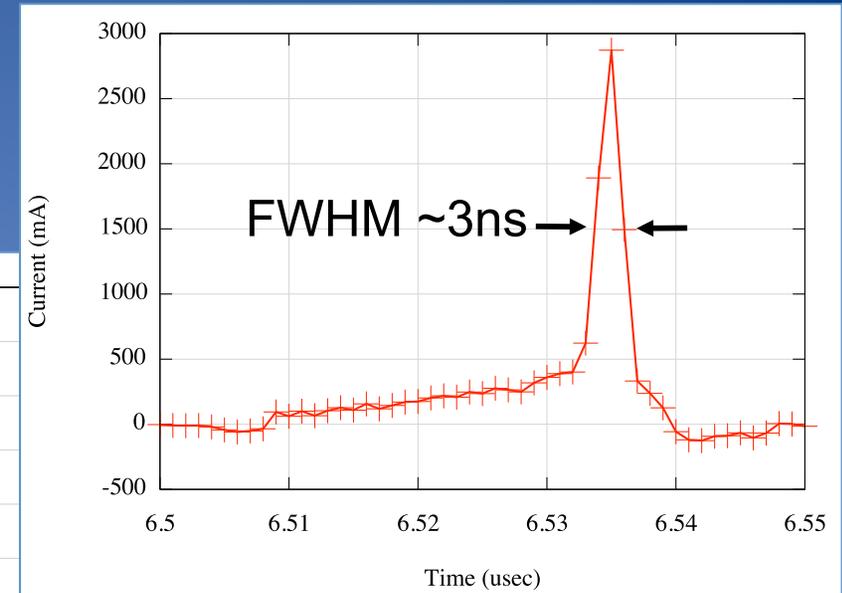
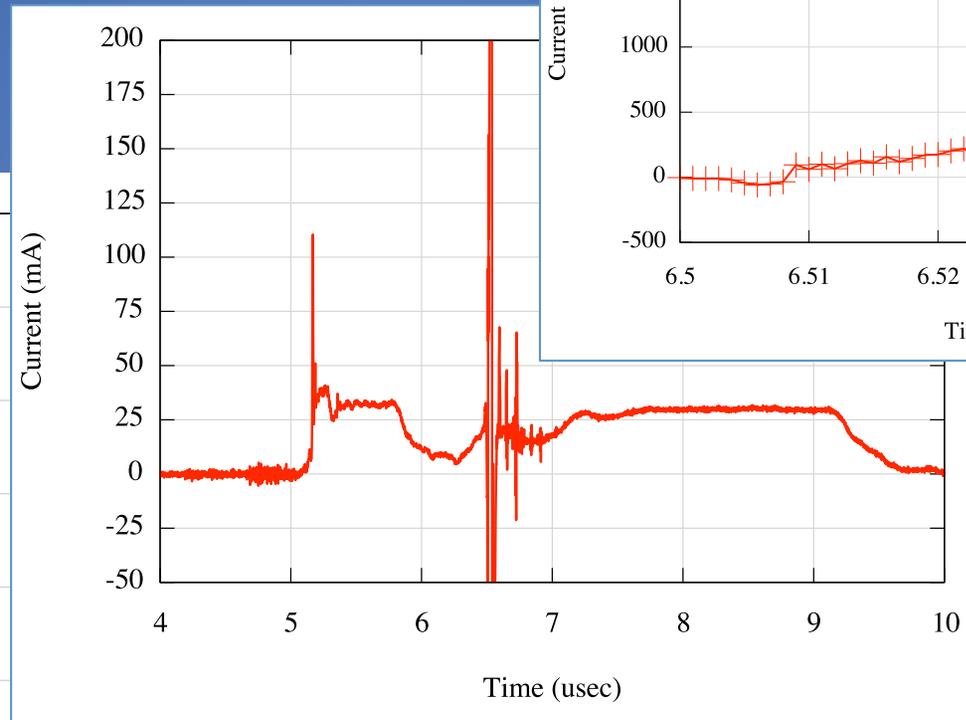
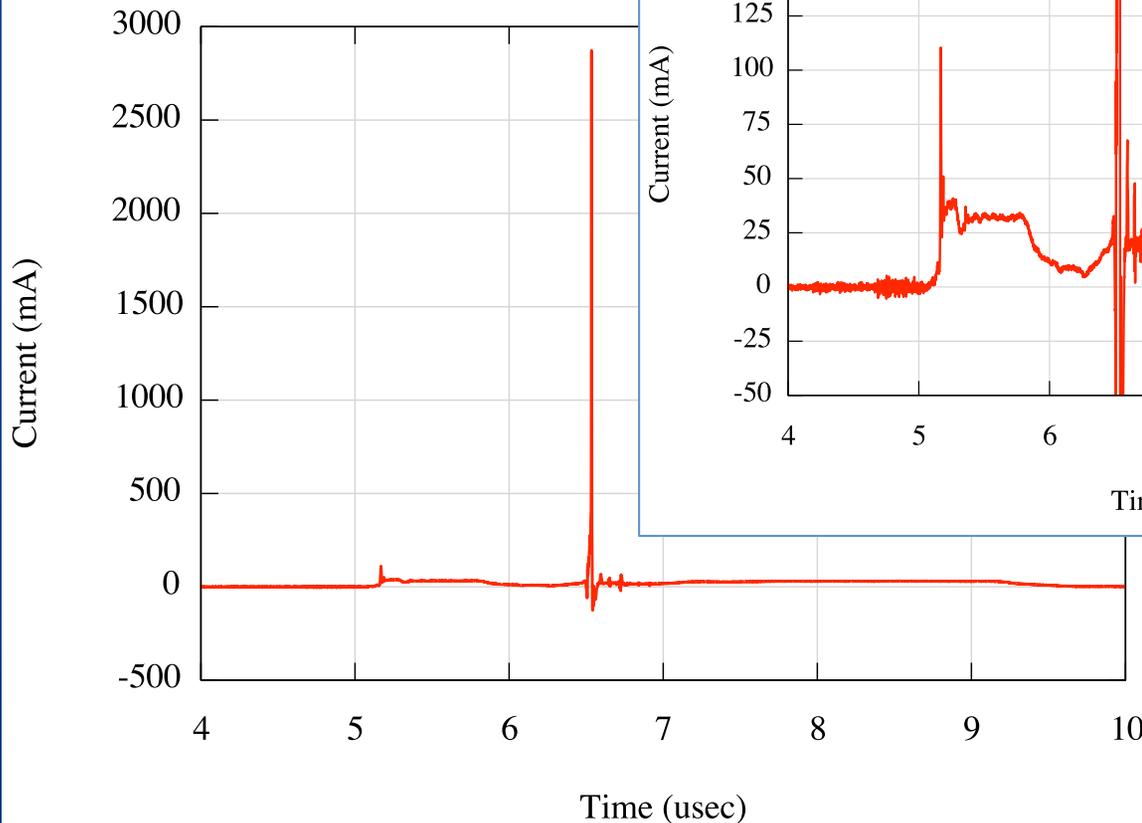


- **Alumina scintillator** with intensified, gated MCP camera
– 10 ns resolution, ~ 30 ns minimum gate

Compressed Current

Peak current $\sim 2.8\text{A}$

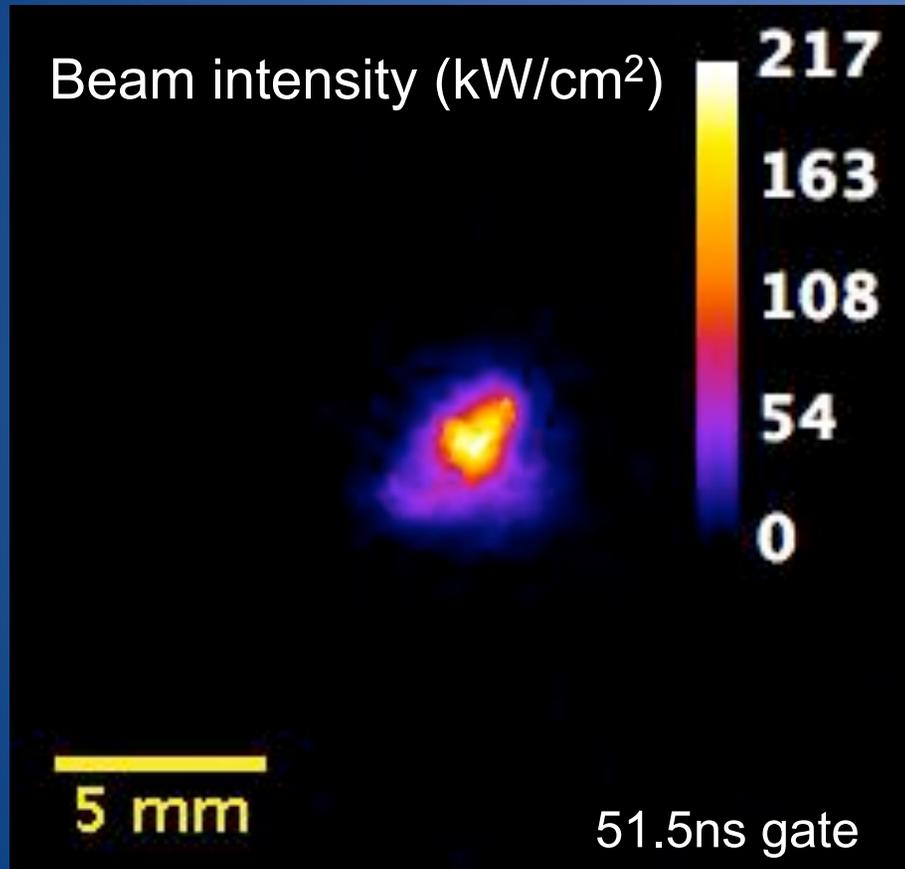
Uncompressed $\sim 30\text{mA}$



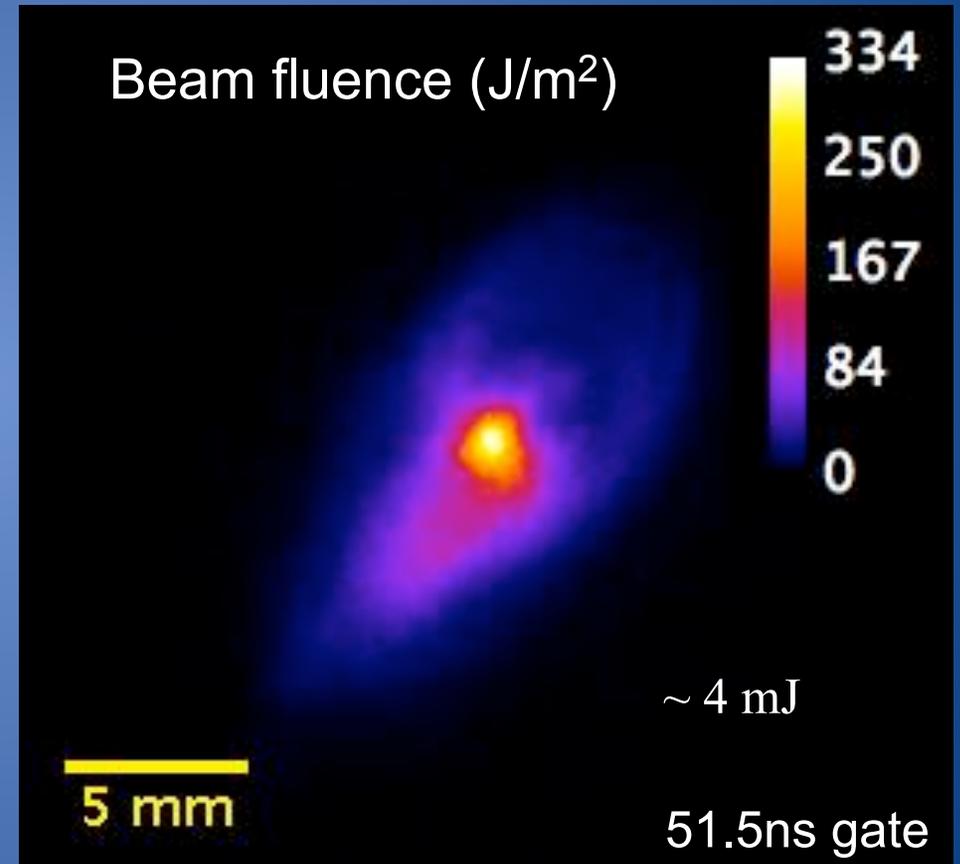
Compression
Ratio >90

Beam Intensity and Fluence

284kV beam energy, 25mA current, ~7kW beam power



Uncompressed

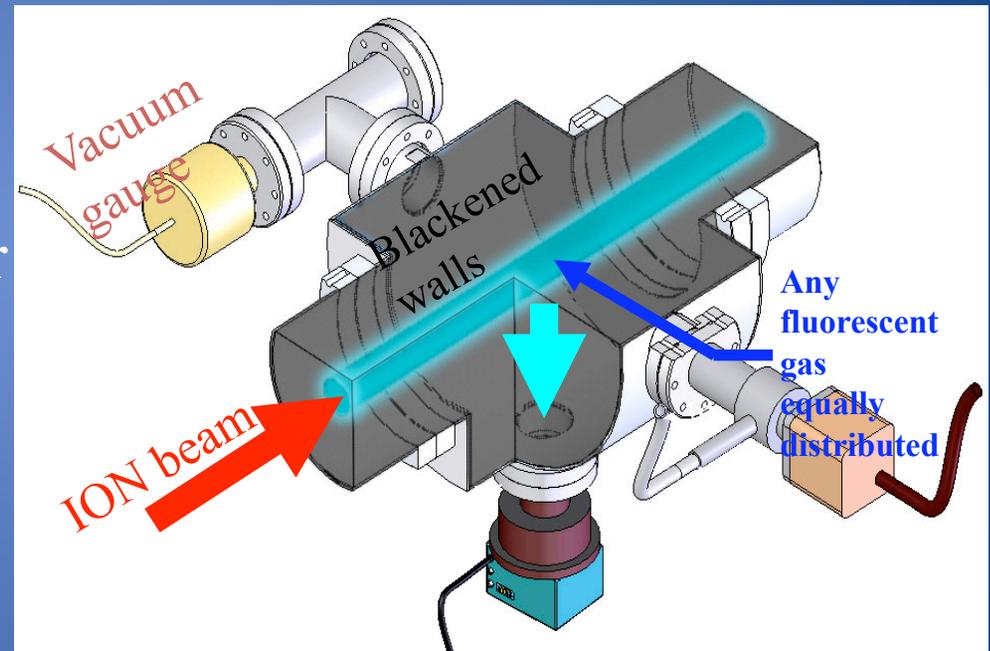


Compressed

High Intensity Beam Profiling

Solid state scintillators incur damage and lose sensitivity with every shot – nonlinear in intensity

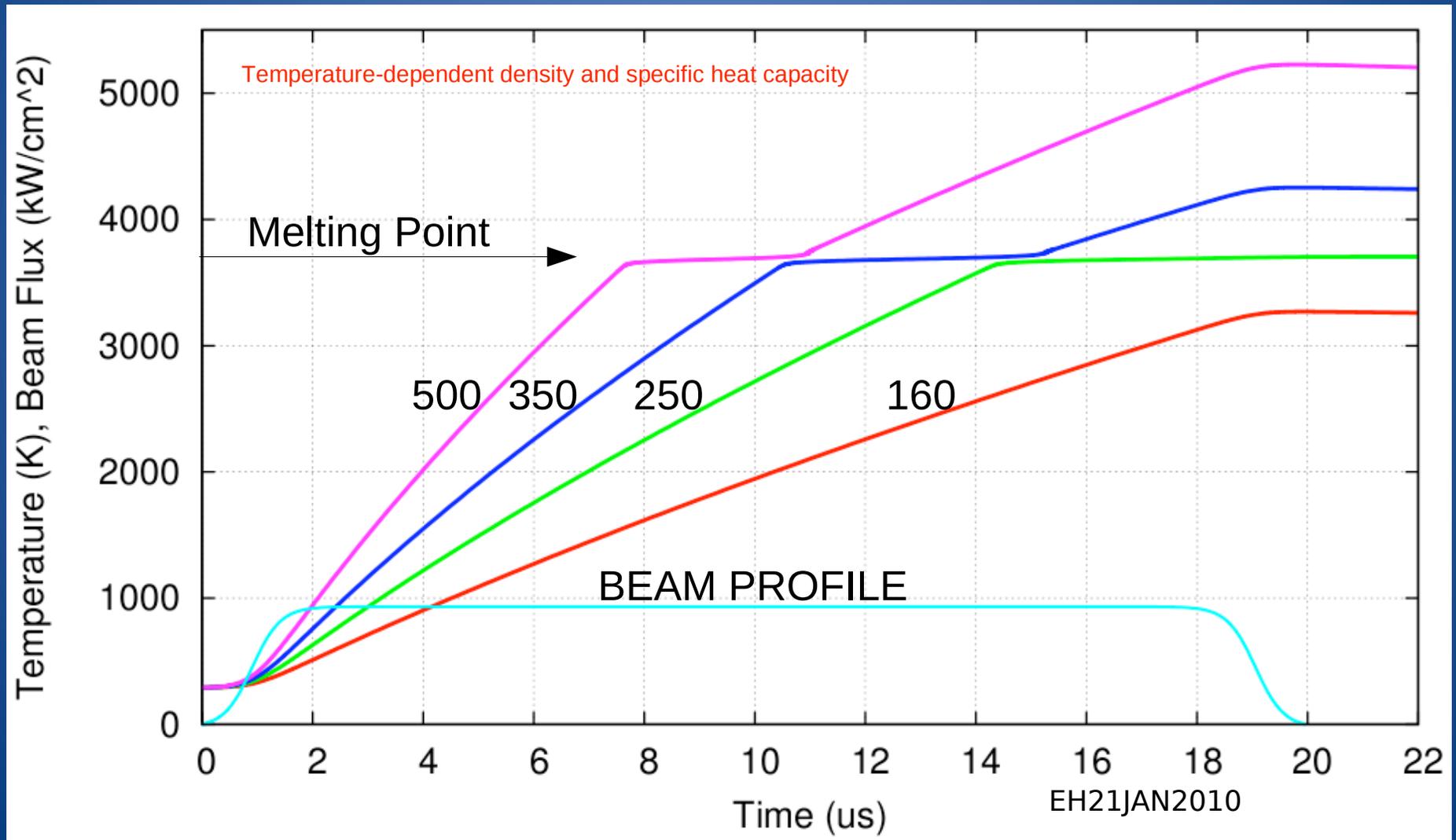
Gas phase Beam Intensity Monitor
Observe beam excited gas lines,
correlate to beam distribution



Tungsten foil calorimeter –
3 μm thick foil responds to beam peak intensity –
radiated intensity $\sim \sigma T^{4+}$

Calibrating the calorimeter

Merging of beam and target diagnostics

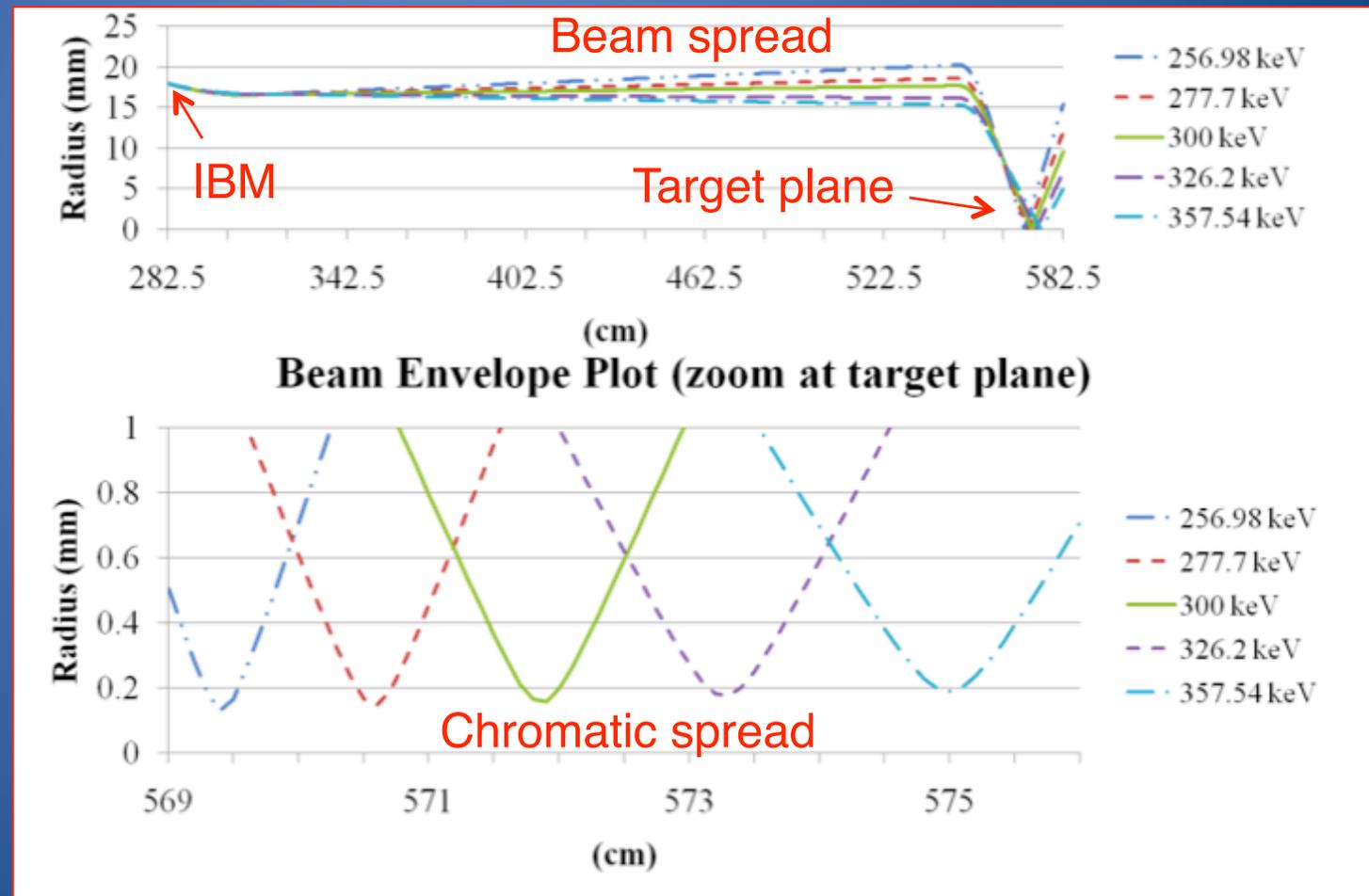


Time-Dependent Aberrations

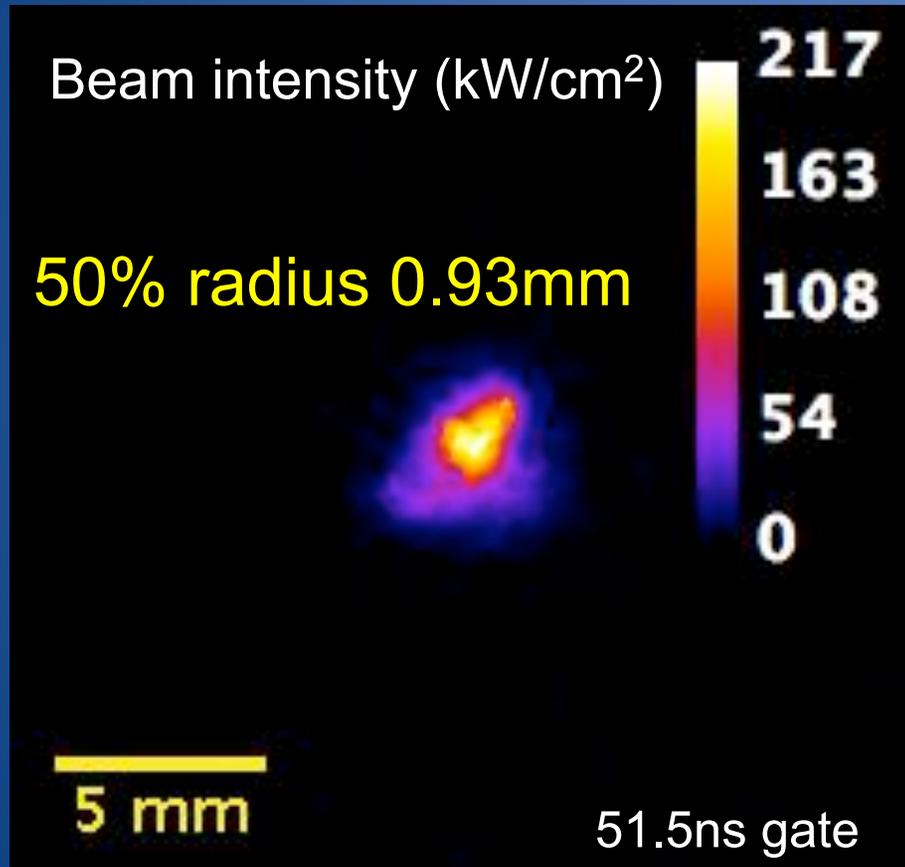
$$\Delta v_r \approx -\frac{qr}{2mv_{z0}^2} \frac{dV_{gap}(t)}{dt}$$

Time-dependent gap defocusing is correlated to slice energy through the impressed IBM energy modulation.

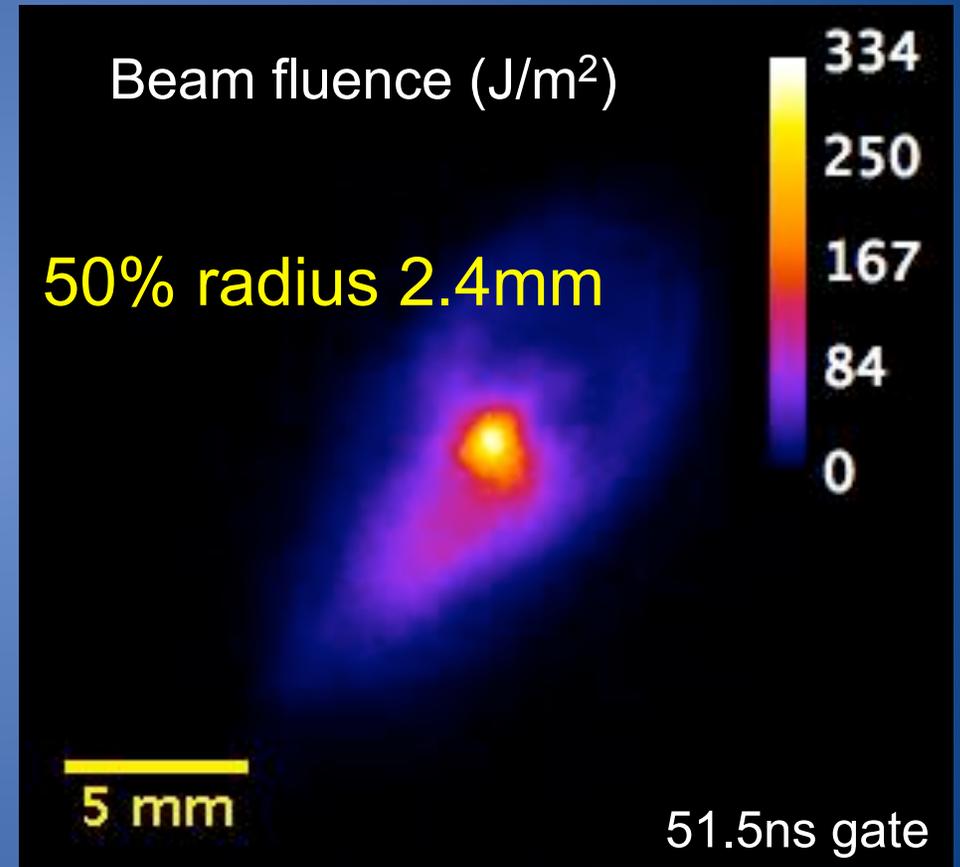
This effect compounds chromatic aberrations in the final focus section.



Measurement of chromatic variations



Uncompressed

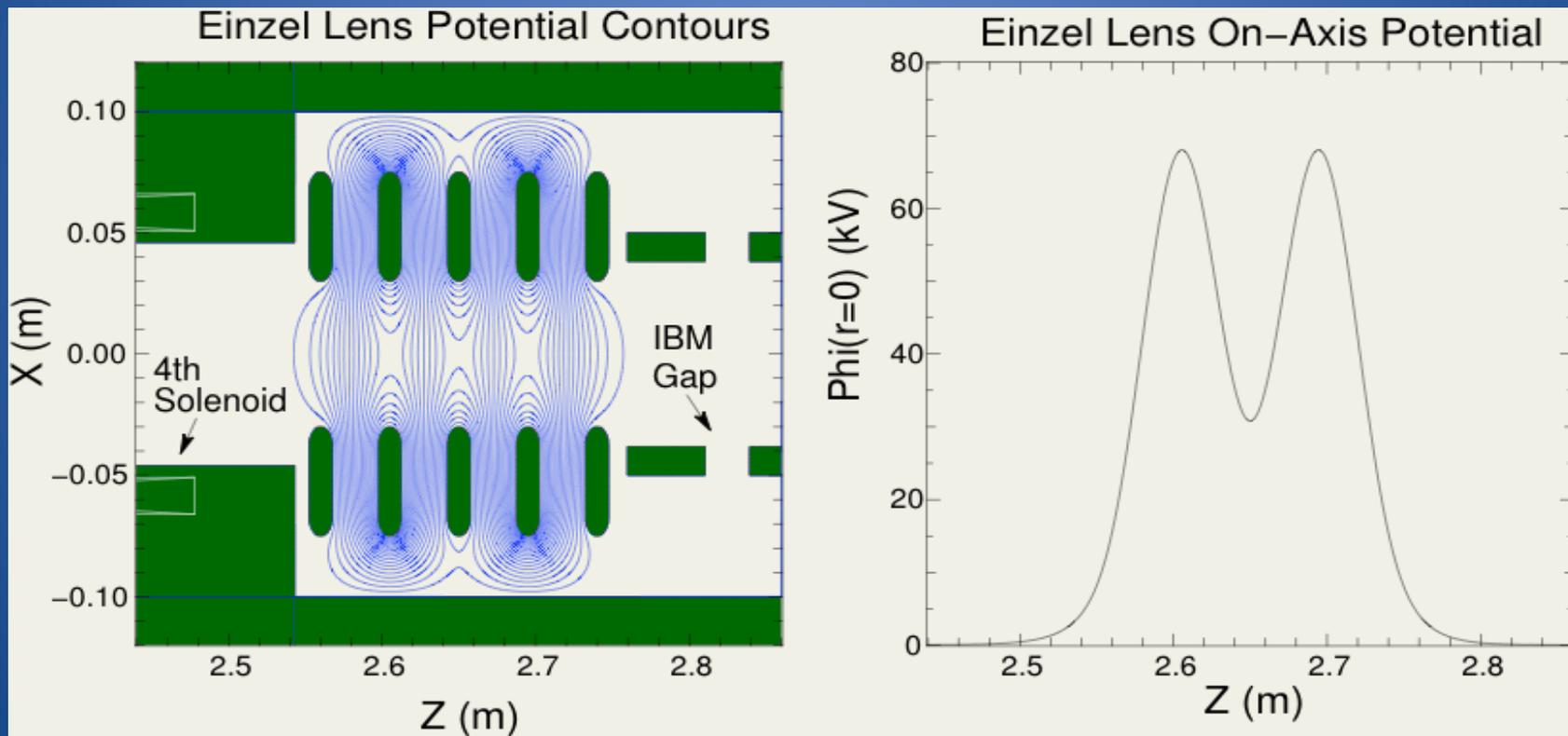


Compressed

~33mJ/cm² compared with ~200mJ/cm²

Fast Corrector

Time-dependent geometric and chromatic aberrations can in principal be corrected.

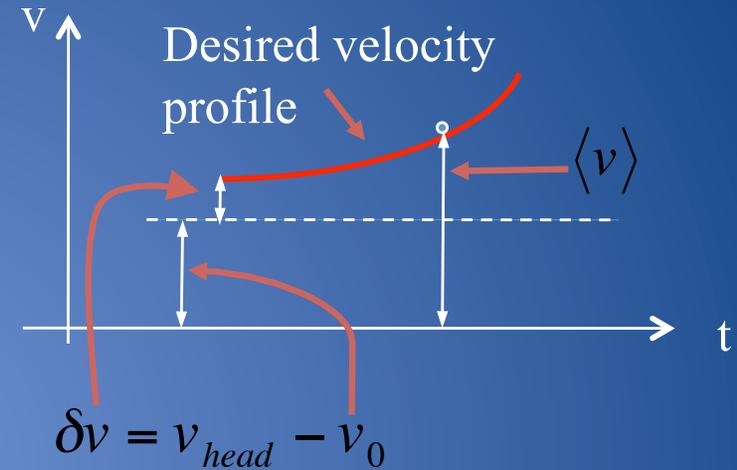


Transit-time effects complicate matters.

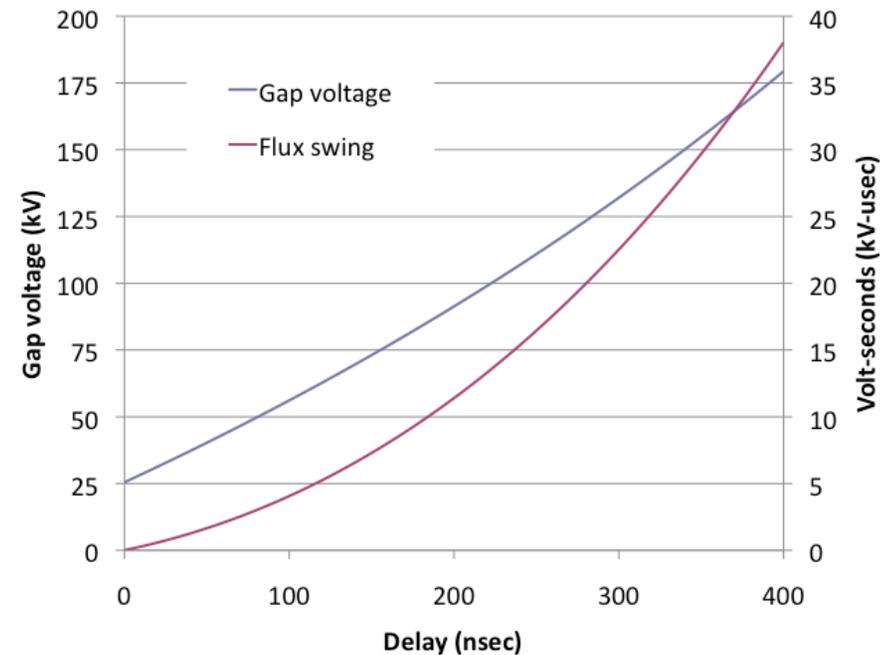
Longitudinal slippage with unipolar IBM waveform

$$\Delta t_{slip}(\Delta z) = \left(\frac{1}{\langle v \rangle} - \frac{1}{v_0} \right) \Delta z$$

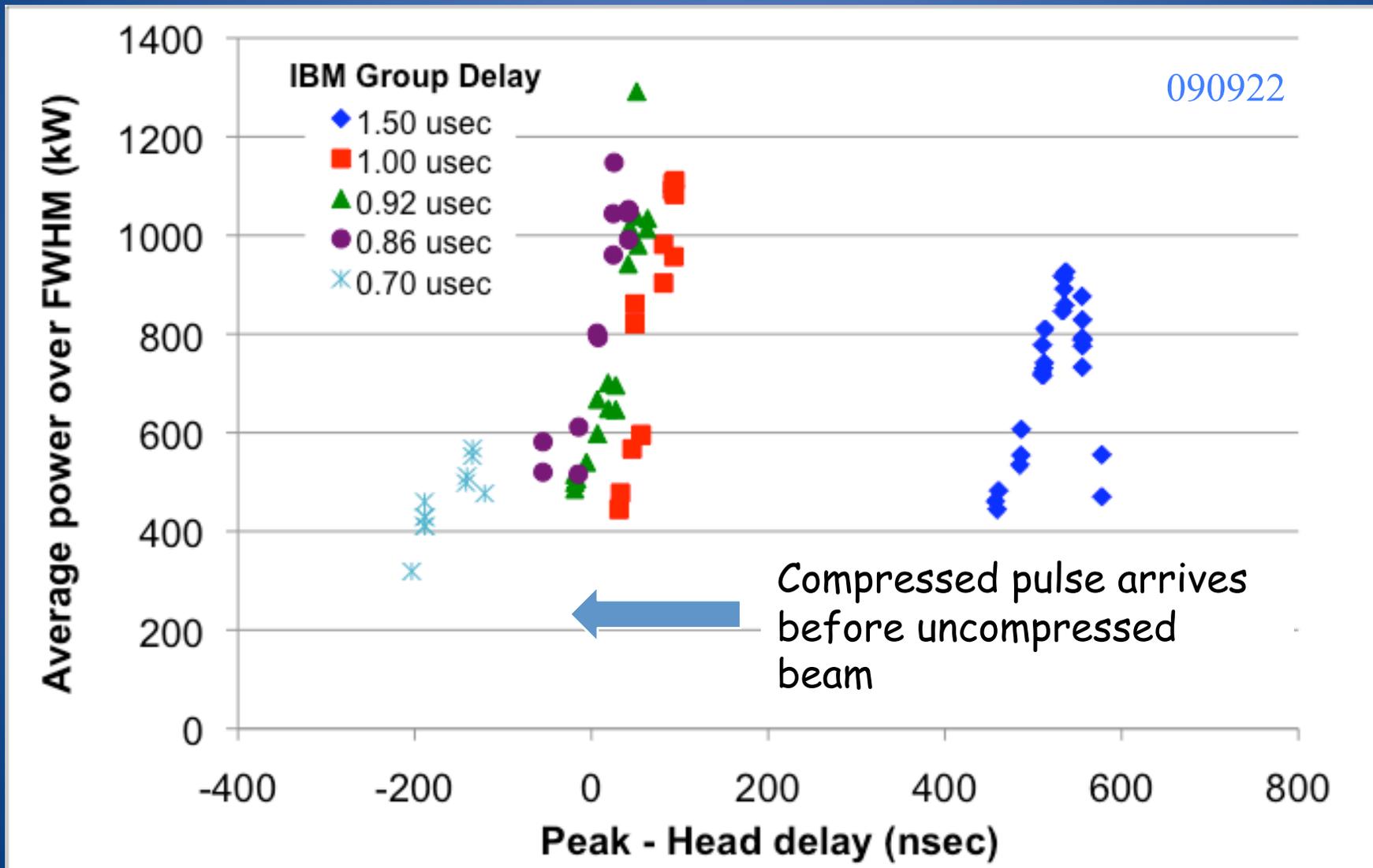
$$\langle v \rangle = \frac{L_{Drift}}{\tau_p} \ln \left(\left[1 - (v_0 + \delta v) \frac{\tau_p}{L_{Drift}} \right]^{-1} \right)$$



IBM and Drift Design			
Ion Energy	E (eV)	300000	
Initial Axial Velocity	v_0 (m/s)	1216819	
Initial Pulse duration	τ (s)	4.00E-07	
Pulse length	$v_0 \cdot \tau$ (cm)	48.7	
Drift length	L_d (m)	2.88	
Avg. Voltage gain	\bar{V} (eV)	95000	
Average beam voltage	V_{beam} (eV)	395000	
Beam head velocity	v_h (m/s)	1267417	
Beam tail velocity	v_t (m/s)	1538184	
Average pulse velocity	$\langle v \rangle$ (m/s)	1396252	
Slippage rate, $1/\langle v \rangle - 1/v_0$	(ns/m)	-106	
Pulse slippage, $L_d(1/\langle v \rangle - 1/v_0)$	(nsec)	-304	(-) earlier, (+) later
Volt-seconds in pulse	D_{phi} (kV-usec)	37.97	400ns
Average voltage across gap	kV	94.9	



Longitudinal compression vs. group delay



Collective focusing

PRL 103, 075003 (2009)

PHYSICAL REVIEW LETTERS

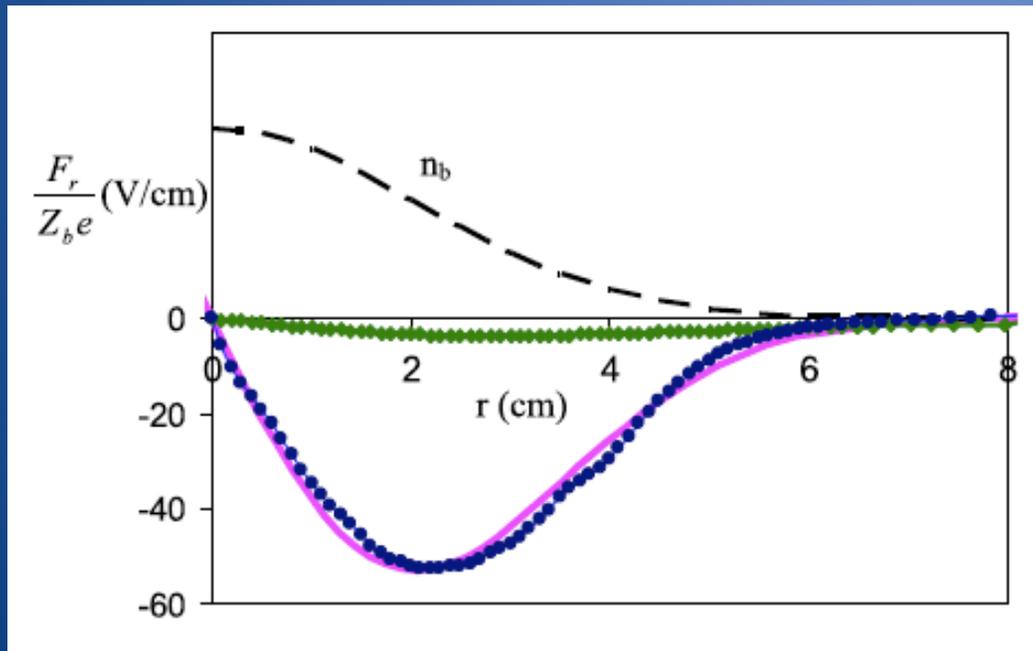
week ending
14 AUGUST 2009

Enhanced Self-Focusing of an Ion Beam Pulse Propagating through a Background Plasma along a Solenoidal Magnetic Field

Mikhail A. Dorf, Igor D. Kaganovich, Edward A. Startsev, and Ronald C. Davidson

Plasma Physics Laboratory, Princeton University, Princeton, New Jersey 08543, USA

(Received 17 February 2009; published 13 August 2009)



A collective focusing scheme has been suggested by colleagues at PPPL.

Requires a final focus magnet field of $\sim 100\text{-}300\text{G}$ rather than 8T !

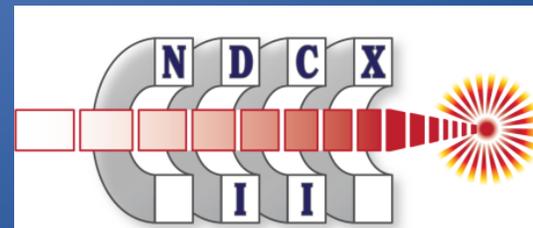
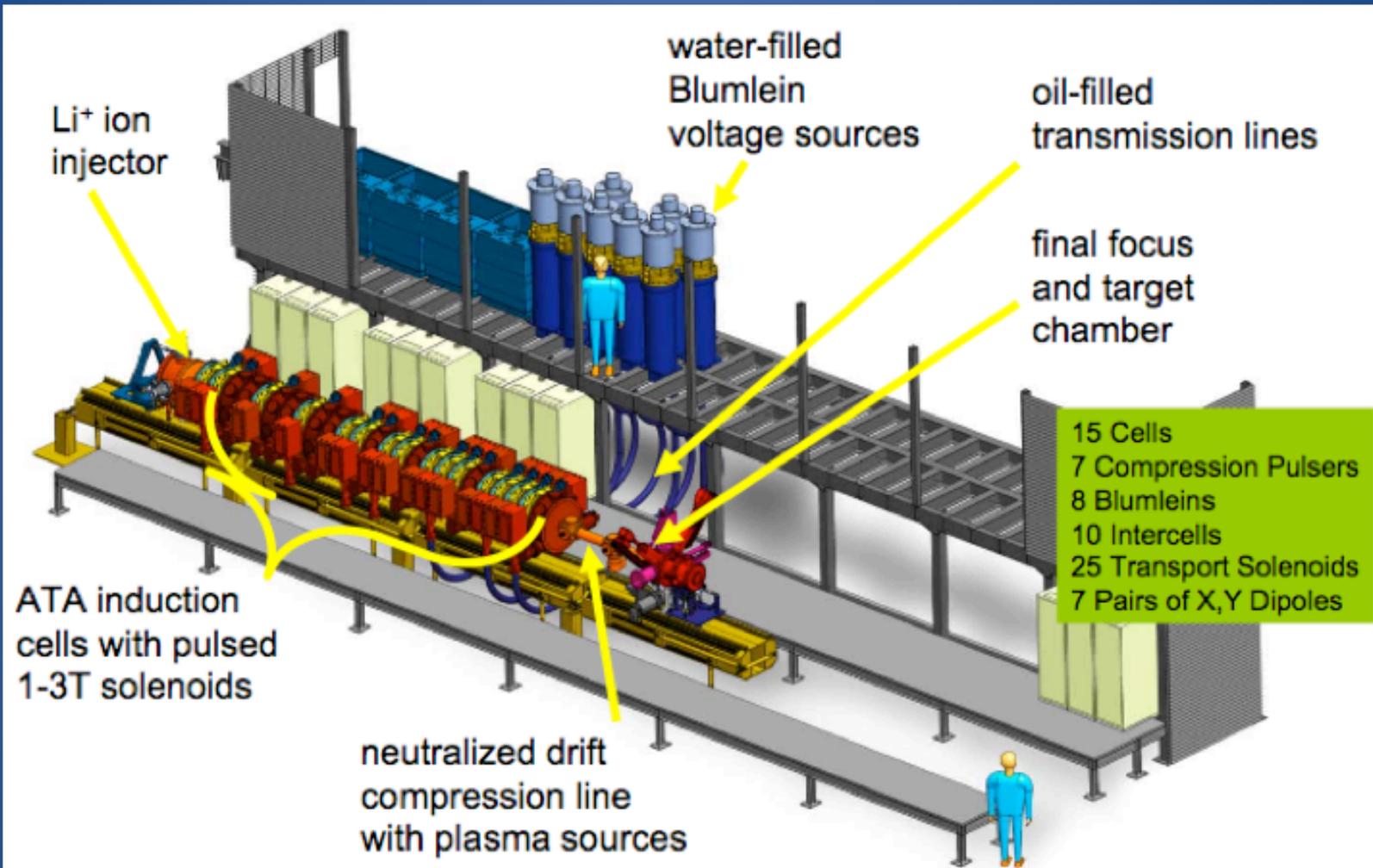
Electrostatic focusing by trapped electron population.

The moving ion beam polarizes the magnetized plasma background, creating a strong radial self-electric field, which provides the enhanced self-focusing.

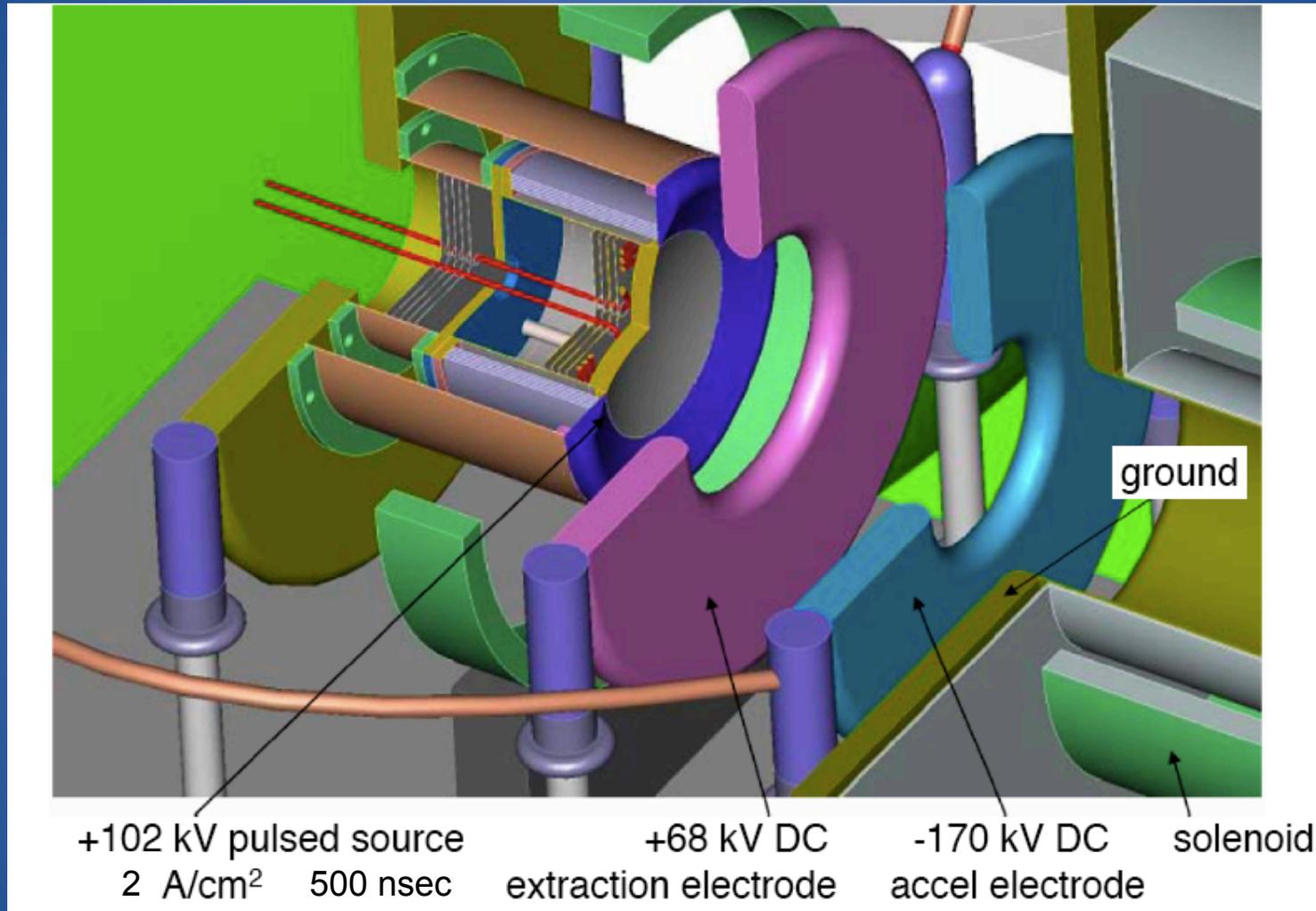
Introduction

- HEDP and WDM
- Why Ion Beams?
- Manipulating Ion Beams for HEDP
 - Stage 1: Neutralized transport and focus (NTX)
 - Stage 2: Neutralized drift compression (NDCX-I)
 - Stage 3: Add acceleration schedule (NDCX-II)
 - Stage 4: Pulse shaping and shock timing (IB-HEDPX)

NDCX-II (15 cells, 2.0MeV)



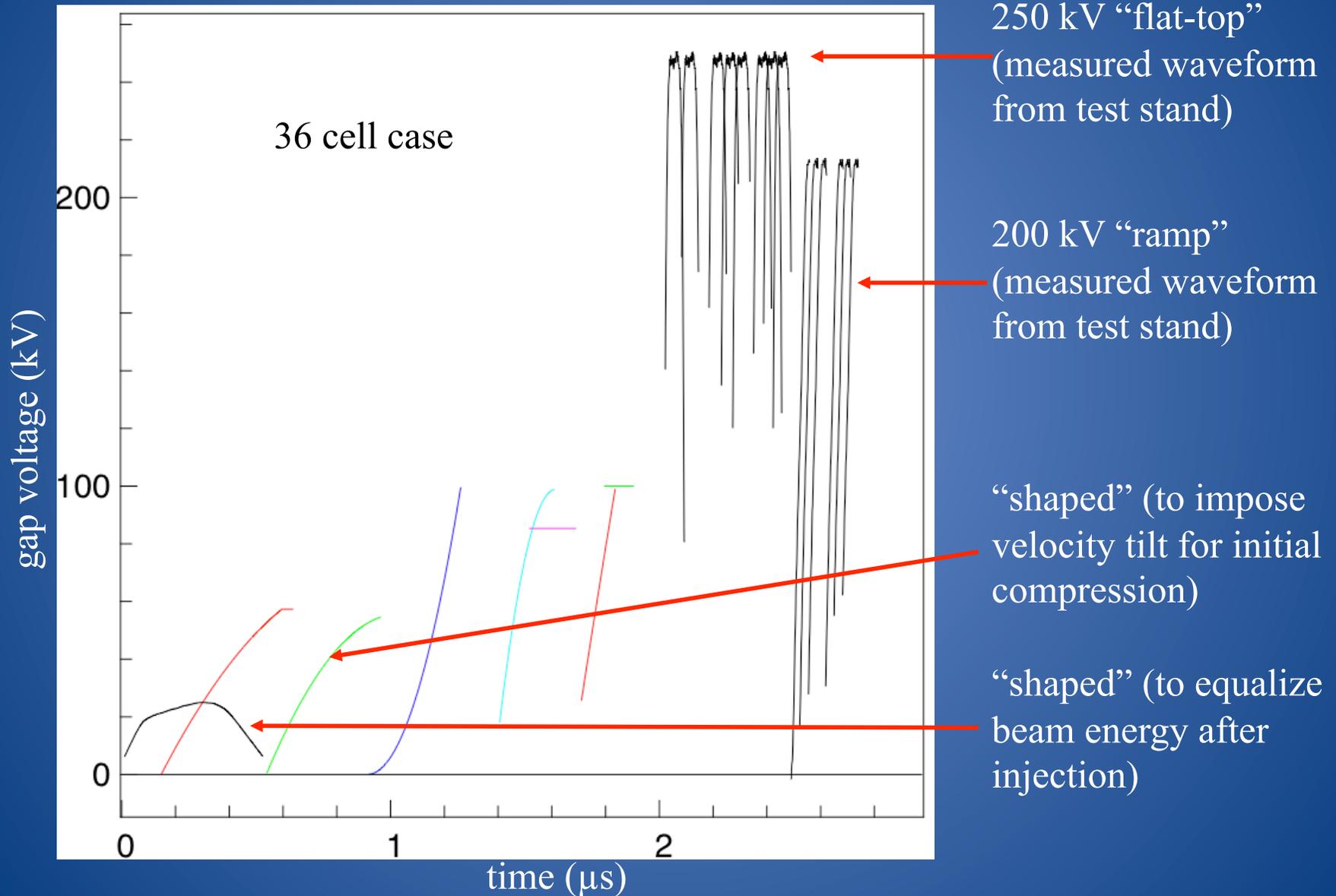
Accel/Decel Injector



Perveance 'multiplier':

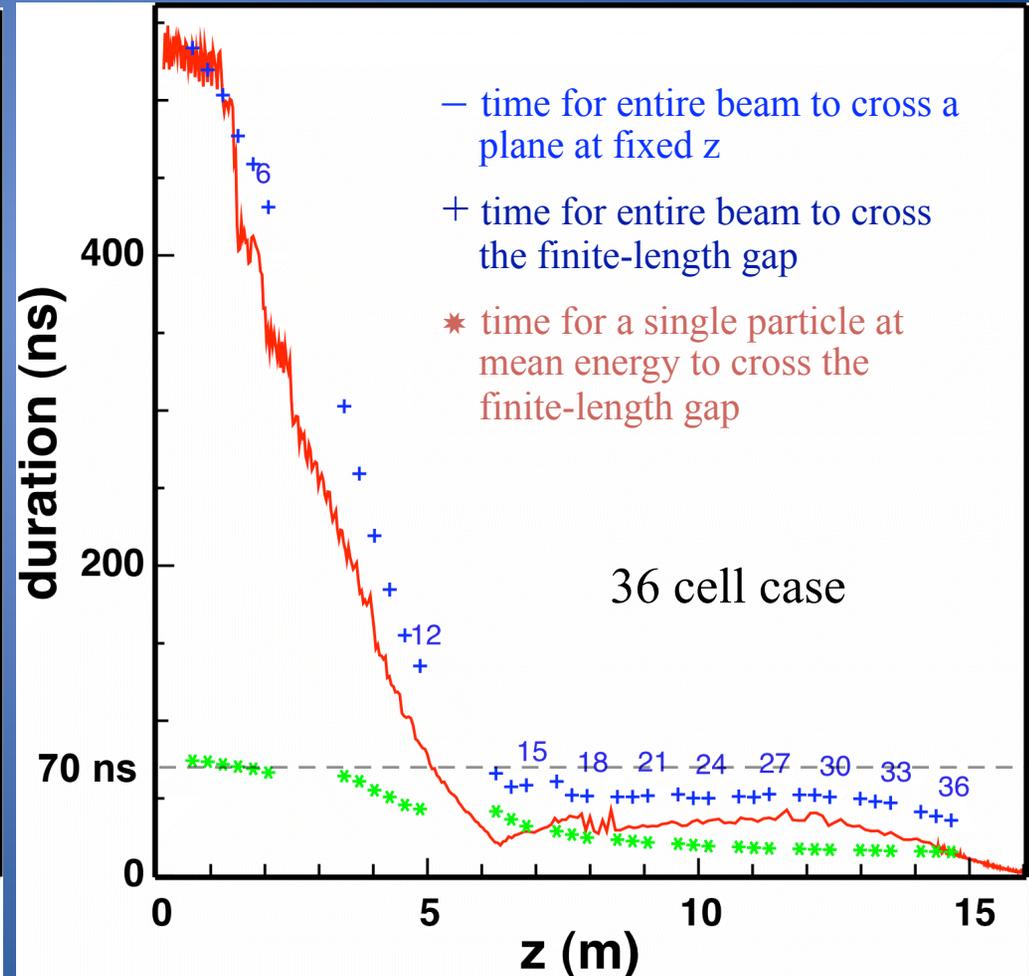
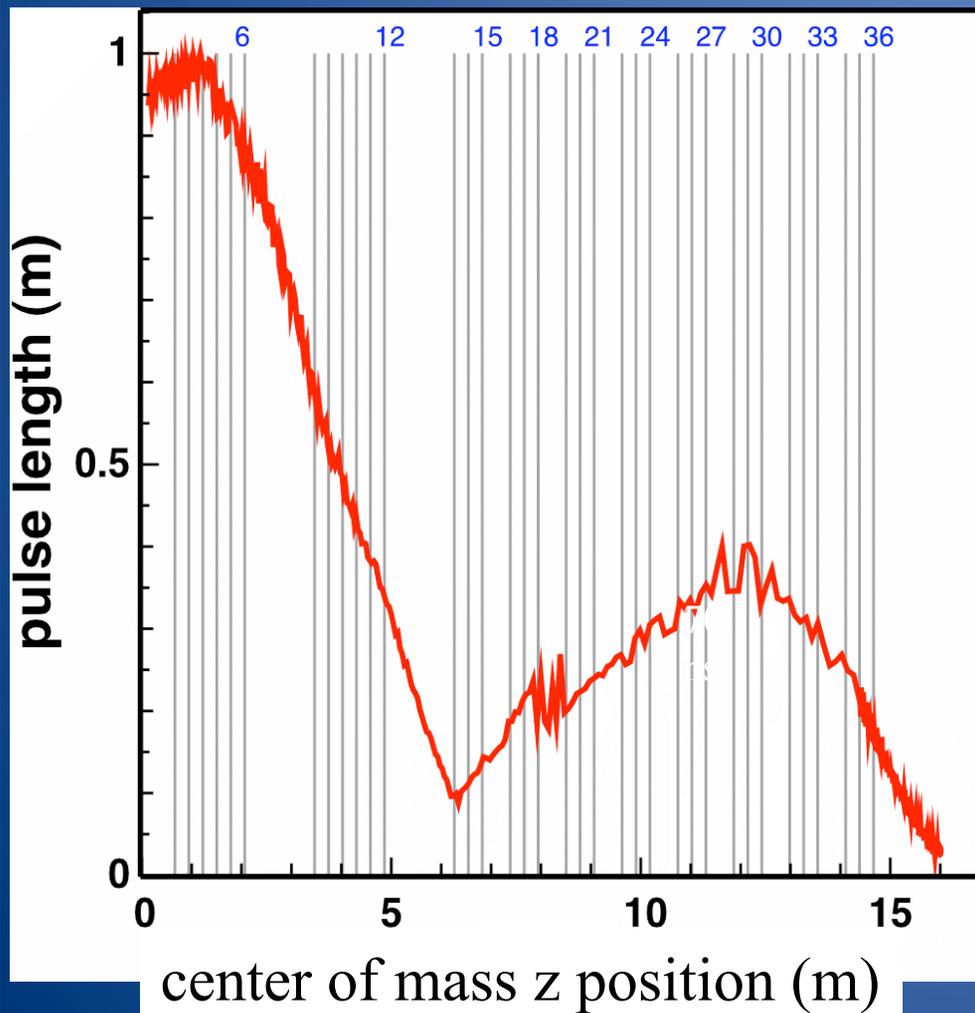
Shortens pulse length to economize on induction module volt-seconds.

Compression and Acceleration



“Let It Bounce”™

Final bunch lengths can be obtained with minimum emittance growth by allowing the longitudinal phase space to evolve, with guidance.

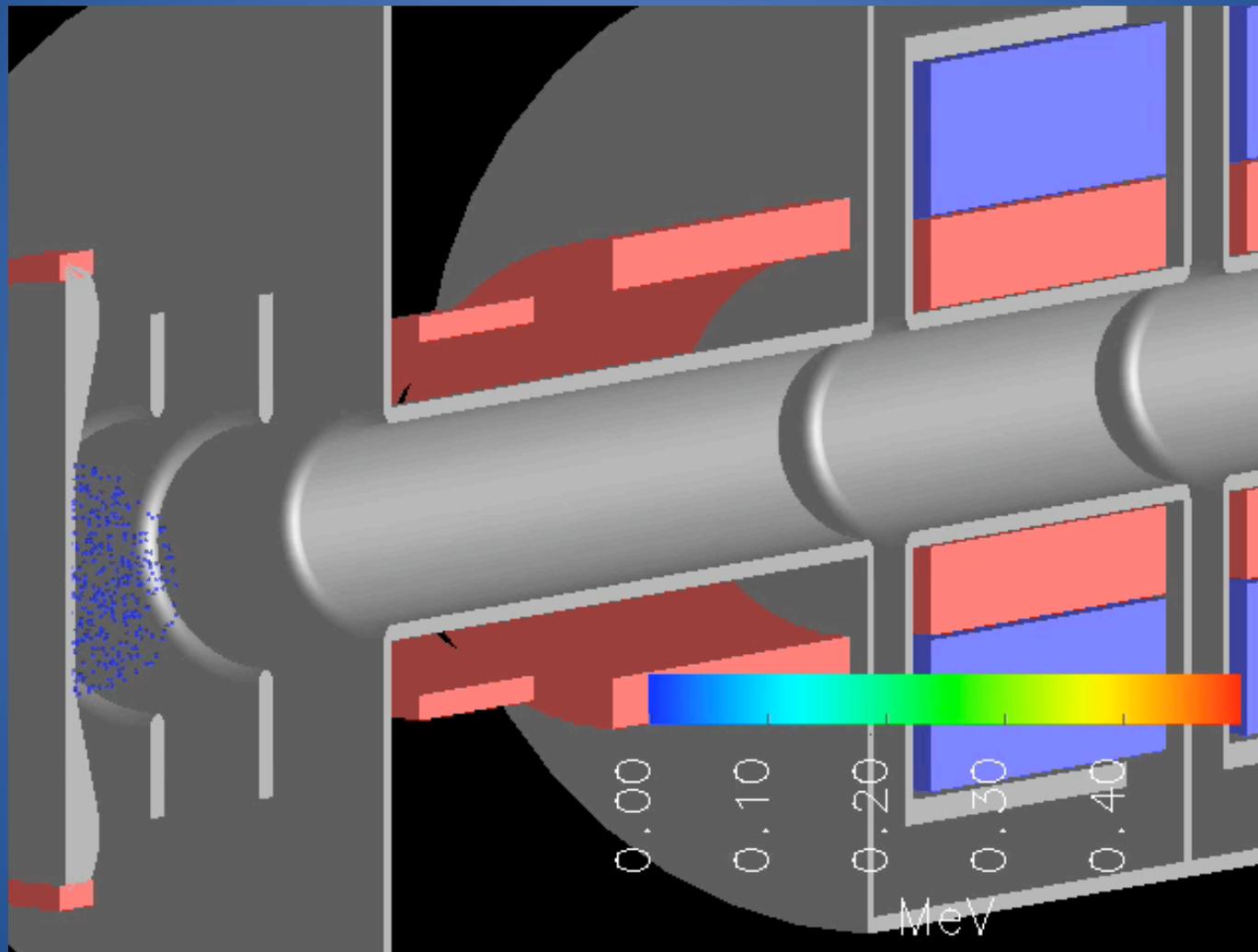


NDCX-II vs. NDCX-I

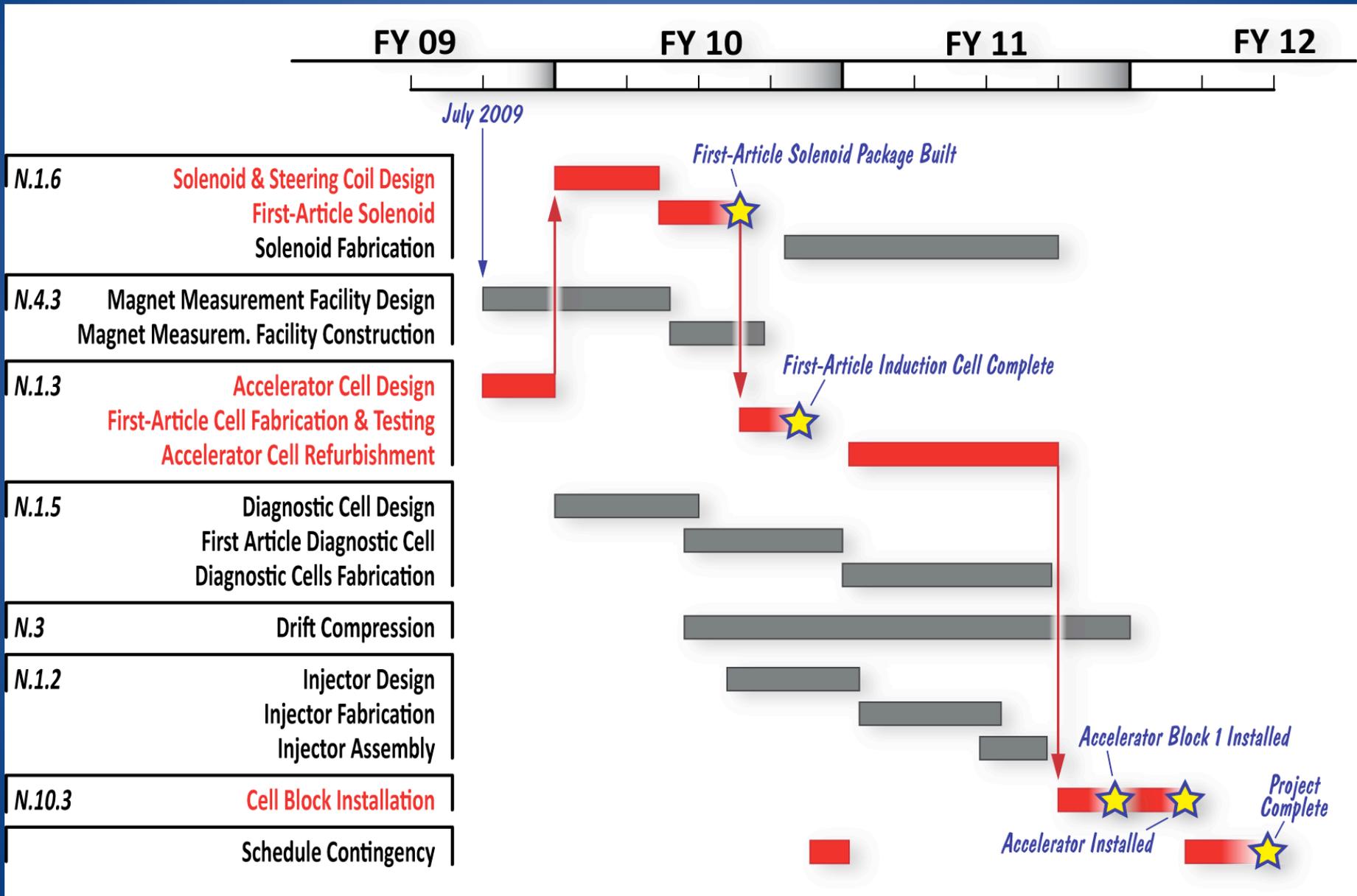
	NDCX-I (typical bunched beam)	NDCX-II 15-cell (r,z simulation)
Ion species	K ⁺ (A=39)	Li ⁺ (A=7)
Total charge	15 nC	55 nC
Ion kinetic energy	0.3 MeV	2.0 MeV
Focal radius (containing 50% of beam)	1 mm	0.7 mm
Bunch duration (FWHM)	2 ns	0.45 ns
Peak current	3 A	17 A
Peak fluence (time integrated)	0.03 J/cm ²	17 J/cm ²
Fluence within a 0.1 mm diameter spot	0.03 J/cm ² (50 ns window)	9.5 J/cm ² (1 ns window)
Fluence within 50% focal radius and FWHM duration (k.e. x I x t / area)	0.014 J/cm ²	1.0 J/cm ²

NDCX-II estimates are from (r,z) Warp runs (no misalignments), and assume no timing or voltage jitter in acceleration pulses, no jitter in solenoid excitation, perfect neutralization, and a uniform non-depleted source; they also assume no fine energy correction (e.g., tuning the final tilt waveforms)

NDCX-II, The Movie



Project Path



Prototype hardware components are tested for QA



IB-HEDPX

- Conceptual device to attain $\sim 10\text{eV}/1\text{Mbar}$ conditions in planar heated targets.
- Extend NDCX-II to 50 ATA cells, $\sim 10\text{ MeV}$ acceleration
- Short pulses (100-200ps)
- Pulse shaping and shock timing
 - Can study hydrodynamic coupling and shock physics.

Summary

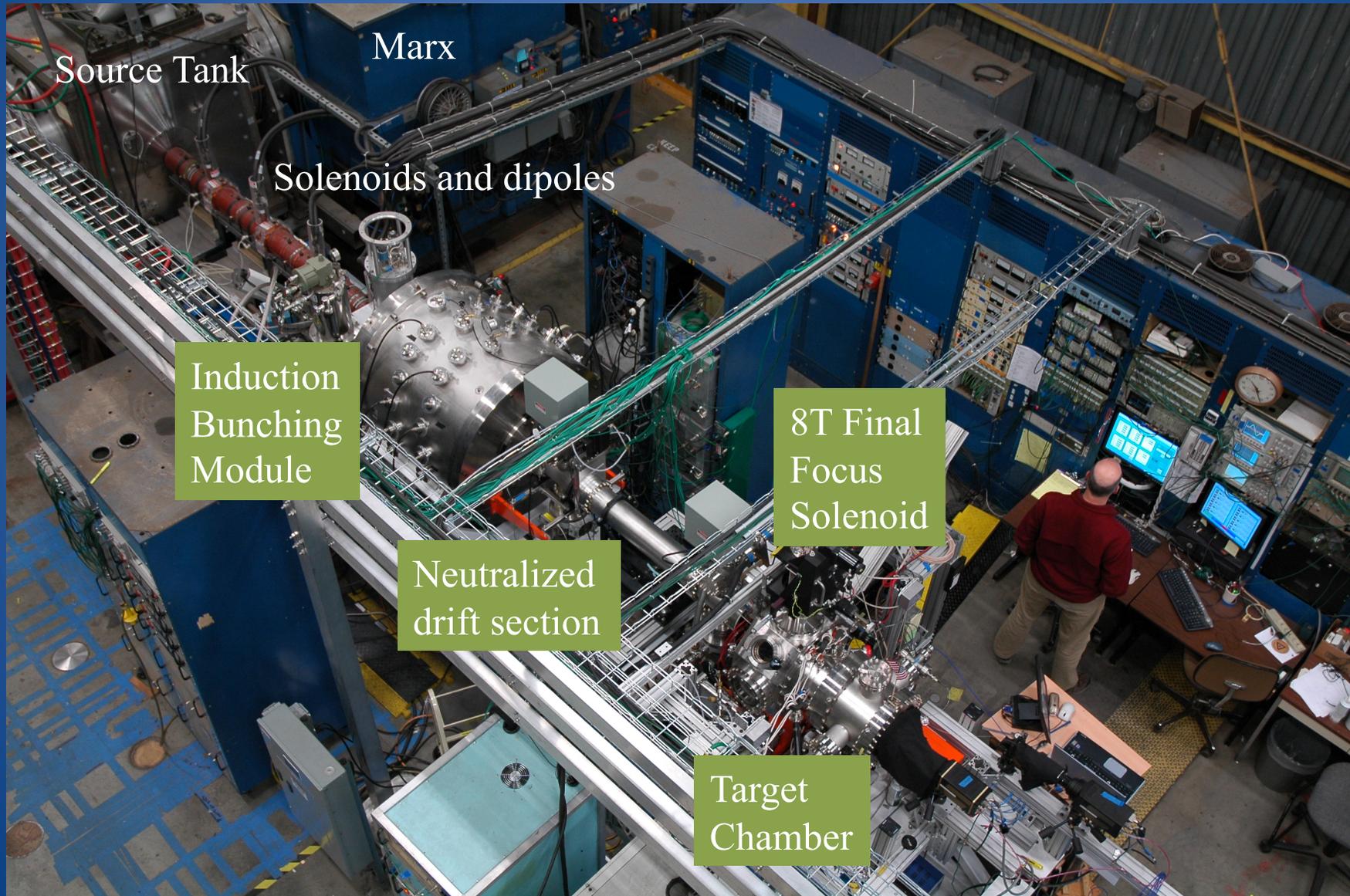
- Beam manipulations are crucial elements in ion beam drivers for HEDP studies.
- These manipulations are only possible due to access to high density, large volume, cool ($T_e \sim eV$'s) plasma sources.
- NDCX-II will significantly expand our reach into WDM and HEDP studies.
- Complex beam dynamics will play an ever increasing role in providing usable beam fluence.

Back up slides

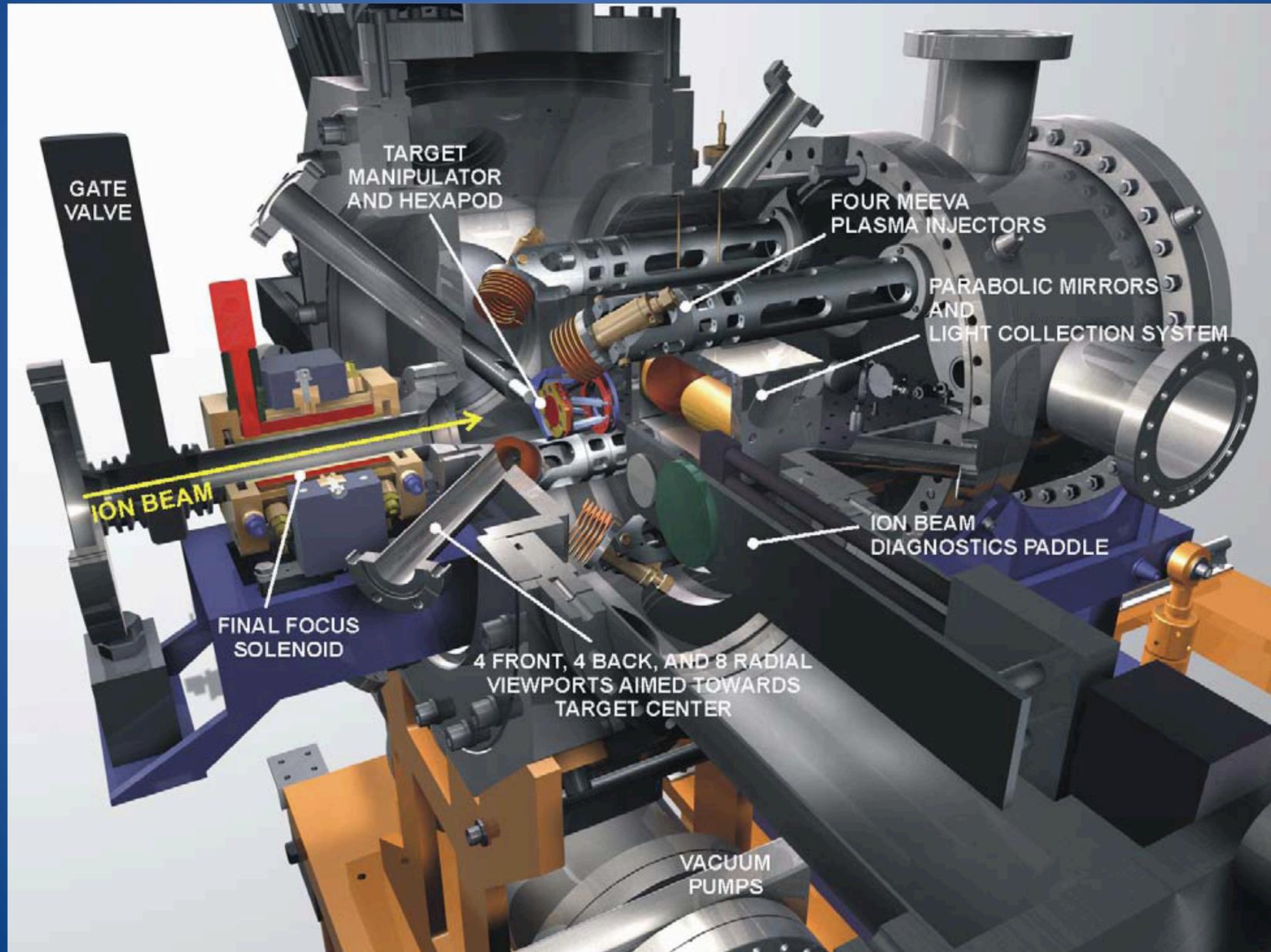
Beam intensity requirements

- Source brightness (space charge limited)
 - Aluminasilicate sources,
 - $K^+, Na^+ \sim 0.5 A / (mm \text{ mrad})^2$ ($J \sim 10 \text{ mA/cm}^2$)
 - $Li^+ \sim 0.005 A / (mm \text{ mrad})^2$ ($J \sim 2 \text{ mA/cm}^2$)
 - Normalized emittance $\sim <1$ to few π mm mrad
- Required target plane peak intensity
 - $T \sim 0.5 \text{ eV}$ (15nC K^+ , 300keV) $0.2 J/cm^2$, $\sim 200 A/cm^2$
 - $T \sim 1-2 \text{ eV}$ (30nC Li^+ , 3MeV) $30 J/cm^2$, $\sim 8 kA/cm^2$
 - $T \sim 10 \text{ eV}$ (30nC Na^+ , 10MeV) $200 J/cm^2$, $\sim 1 MA/cm^2$

NDCX-I in Bldg 58



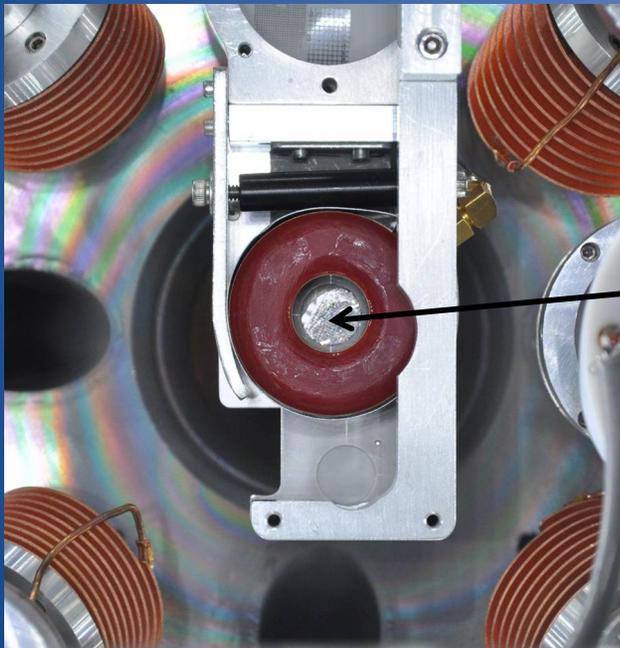
Target Chamber



Target Diagnostics

Fast Optical Pyrometer

Fast Current Transformer



Target Foil

Absolutely-calibrated streak-spectrometer

Challenges and requirements:

- Being able to detect at least ~1 mW of light from 400 nm to 2000 nm in a sub-ns time scale:
- Ultra fast response (sub-ns)
- Higher sensitivity (≥ 1500 K)
- Different, more efficient beam splitting mechanism
- No published paper

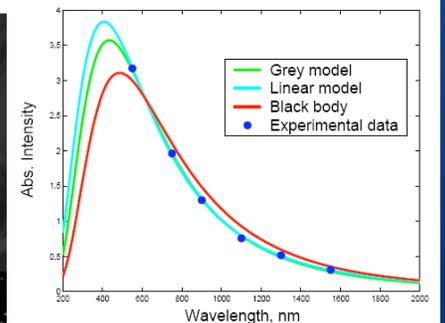
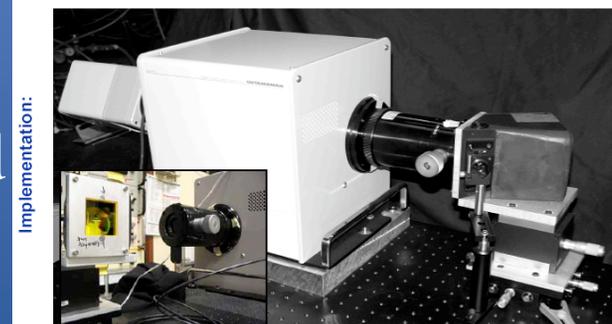
Technical issues:

- Need low noise amplified photo receiver with flat gain from DC to 4GHz and big active surface.
- Fiber coupling limits the efficiency
- Picoseconds time scale means: Modal (temporal) dispersion in MM fiber, careful cabling, impedance match and termination of detectors

Channel #1: 750 nm \pm 75 nm, 76 ps rise/fall time detect: level ~2500 K (blackbody)
 Channel #2: 1000 nm \pm 75 nm, 72 ps rise/fall time detectable level ~1500 K (blackbody)
 Channel #3: 1400 nm \pm 75 nm, 70 ps rise/fall time detectable level ~2000 K (blackbody)

Working scheme:

- Records continuous spectrum from 500 nm to 850 nm
- Temporal resolution down to 5 ps
- Linear response
- When calibrated with tungsten lamp, can be used for true temperature determination and spectral emissivity
- Can discriminate useful signal from NDCX plasma radiation
- No published paper



NDCX-II Induction Cell Detailed Design

A module is one or more induction cells plus an “Intercell” for beam steering and diagnostics

